

Lecture at Xidian University
On frontiers of modern optics

Scattering of shaped beam by particles and its applications

IV. Ray theory of wave and its applications

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西安电子科技大学
现代光学前沿专题

波束散射理论和应用

第四讲：波的射线理论及其应用

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Plan of lecture

I. Fundamentals of geometrical optics

II. Geometrical optics for light scattering

III. VCRM - Vectorial complex ray Model

Fundamentals of geometrical optics

What is it used for ?

➤ Imaging in daily life:

- Image by reflection: mirror, water surface, summer route ...
- Image by refraction: fish, stick in water, mirage in the nature ...

➤ Optical instruments:

- Camera, telescope, microscope, ..
- Industrial measurement systems,
- Optical fiber for telecommunication

➤ Scientific research:

- Fluid mechanics: PIV, holography, LDV, PDA,
- Measurement of temperature, size distribution, ...
- Biological imaging, ...

Advantages:

- Simple
- Object of arbitrary shape

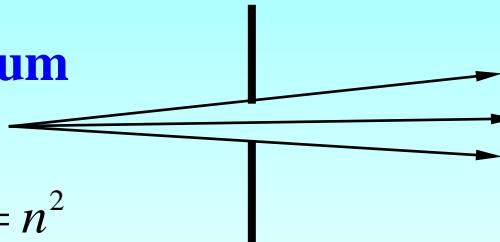
Fundamentals of geometrical optics

Condition of application

$$\lambda \ll l$$

Wavelength λ much smaller than the dimension of the object l .

1. Straight line in homogeneous medium



2. Eikonal eq. (程函方程): $(\nabla S)^2 = n^2$

3. Eq. of ray: $\nabla n = \frac{d}{ds} \left(n \frac{d\vec{r}}{ds} \right)$

The rays are perpendicular to the wave front.

Fundamentals of geometrical optics

Reflection and refraction laws

Reflection and refraction on a surface between two media

Law of Snell-Descartes

Reflection:

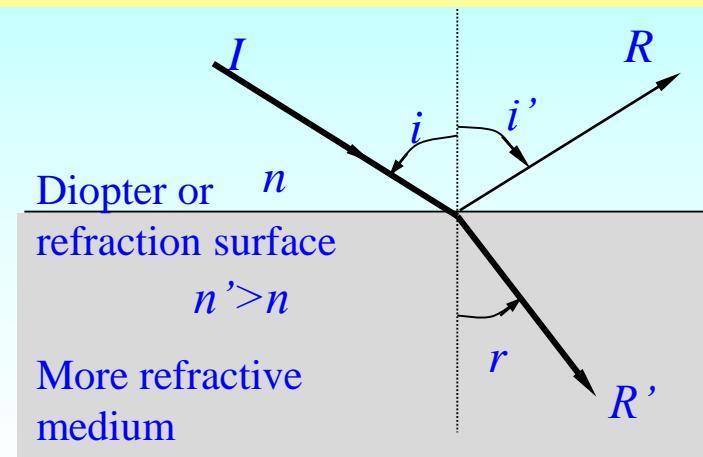
$$i = i'$$

Refraction:

$$n \sin i = n' \sin r$$

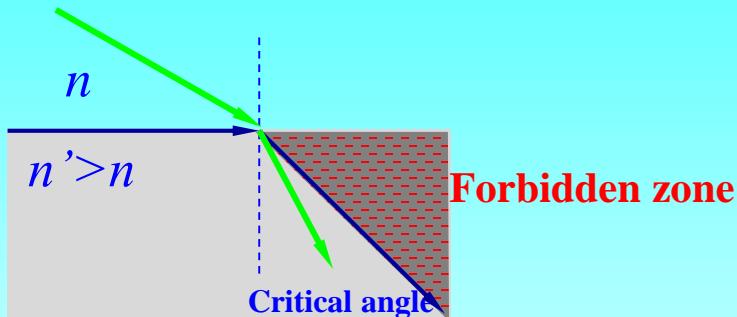
The velocity of light v depends on the refractive index n of the medium:

$$n = \frac{c}{v}$$



Fundamentals of geometrical optics

Total reflection

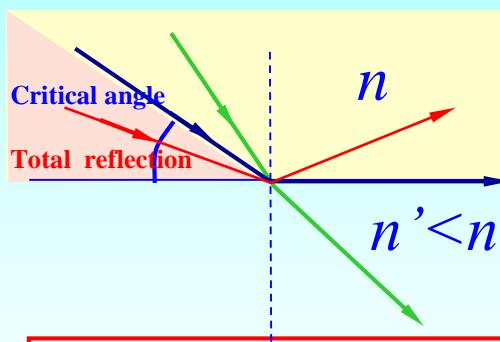


Critical angle:

$$i_l = \arcsin\left(\frac{n'}{n}\right)$$

Total reflection :

$$i \geq i_l$$



Examples :

$$n_2 / n_1 = 1 / 1.333 \quad i_l = 48.6^\circ$$

$$n_2 / n_1 = 1 / 1.500 \quad i_l = 41.8^\circ$$

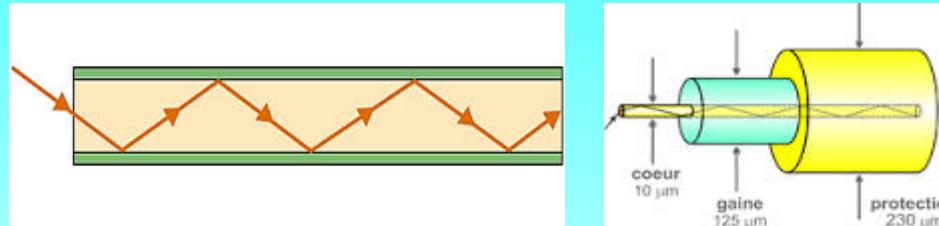
Applications

- Optical fiber
 - Telecommunication
 - Transport de laser beam
- Critical angle for measuring bubbles
- Optical gauge
- Natural mirage

Fundamentals of geometrical optics

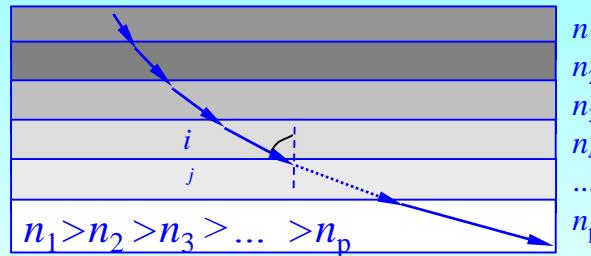
Total reflection - applications

Application 1:
Optical fiber



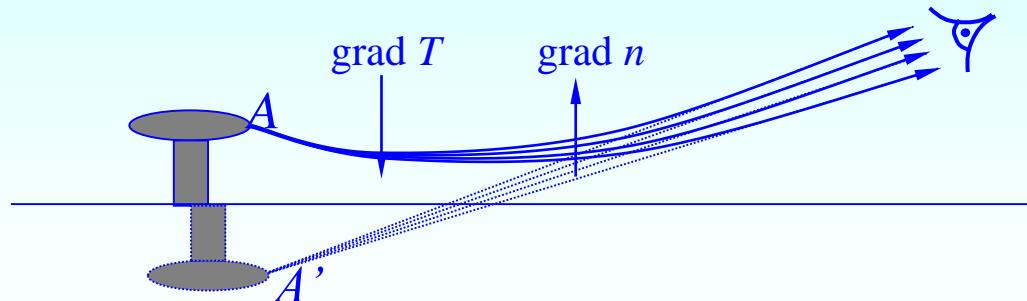
Application 2:
Effect of mirage

Propagation of a light beam
in a parallel stratified
environment:



$$n_j \sin i_j = C$$

$n_i \downarrow$ until $i_j = 90^\circ$,
i.e. total reflection.



Fundamentals of geometrical optics

Relation of amplitudes and intensities

Fresnel formulas – relation between the amplitudes of the reflected/refracted and incident waves

$$r_X = \frac{E_X^r}{E_X^i}, \quad t_X = \frac{E_X^t}{E_X^i}$$

$X = \perp$ ou \parallel polarisation state

E_X^i : amplitude of incident wave

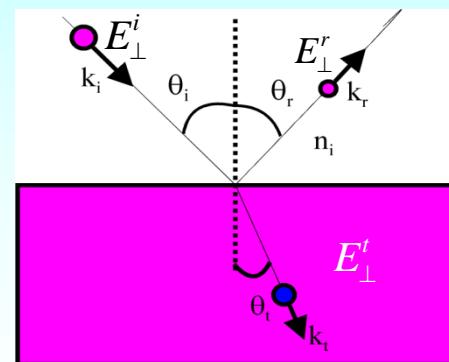
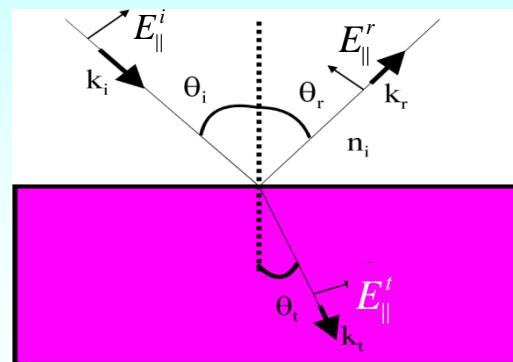
E_X^r, E_X^t amplitudes of reflected/refracted waves

$$r_{\parallel} \equiv \frac{n_i \cos \theta_i - n_r \cos \theta_r}{n_i \cos \theta_r + n_r \cos \theta_i} = \frac{\tan(\theta_i - \theta_r)}{\tan(\theta_i + \theta_r)}$$

$$r_{\perp} \equiv \frac{n_i \cos \theta_i - n_r \cos \theta_r}{n_i \cos \theta_r + n_r \cos \theta_i} = -\frac{\sin(\theta_i - \theta_r)}{\sin(\theta_i + \theta_r)}$$

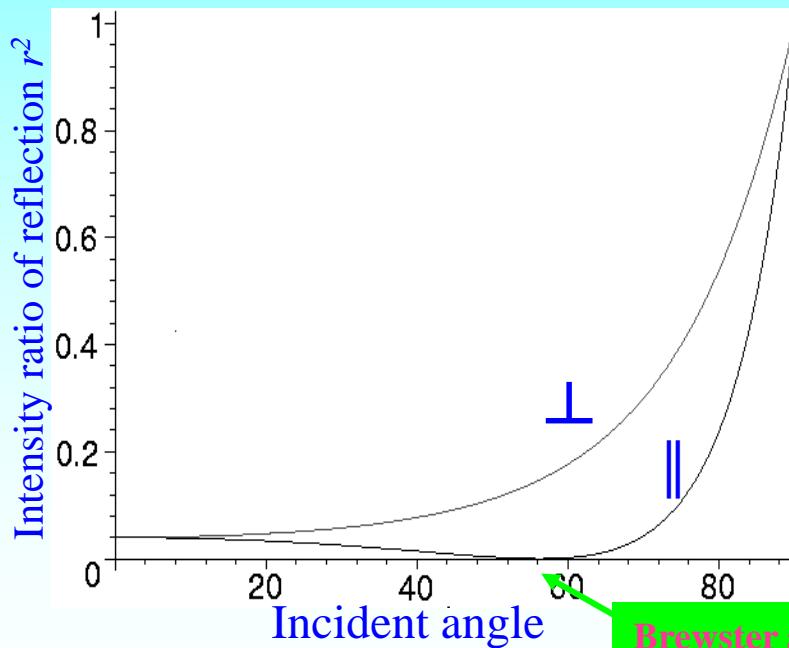
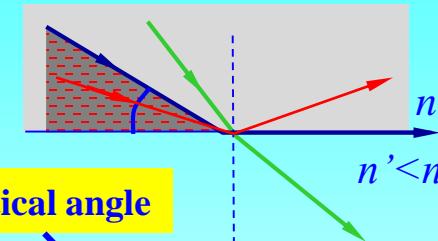
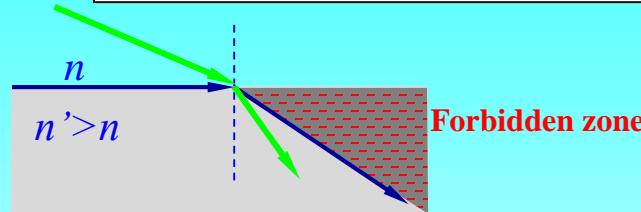
$$t_{\parallel} \equiv \frac{2n_i \cos \theta_i}{n_i \cos \theta_r + n_r \cos \theta_i} = \frac{2 \sin \theta_r \cos \theta_i}{\sin(\theta_i + \theta_r) \cos(\theta_i - \theta_r)}$$

$$t_{\perp} \equiv \frac{2n_i \cos \theta_i}{n_i \cos \theta_r + n_r \cos \theta_i} = \frac{2 \sin \theta_r \cos \theta_i}{\sin(\theta_i + \theta_r)}$$

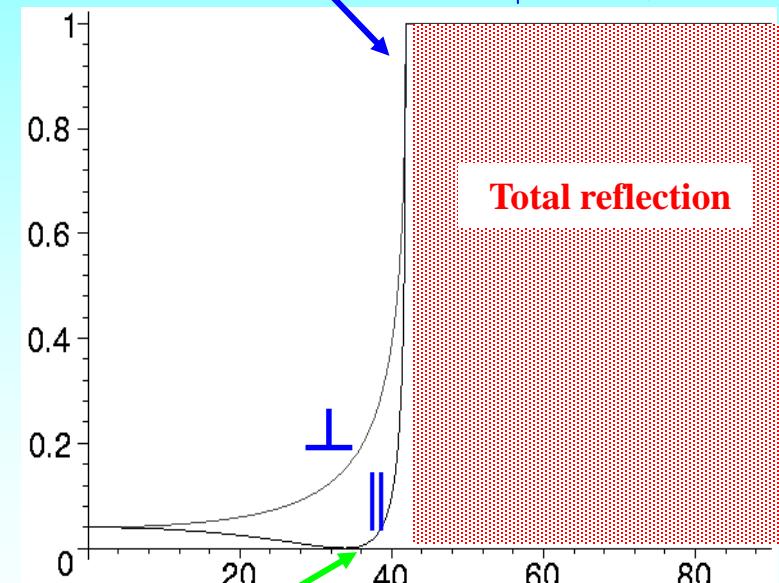


Fundamentals of geometrical optics

Relation of amplitudes and intensities



Brewster angle: $\theta_i + \theta_r = \pi/2$



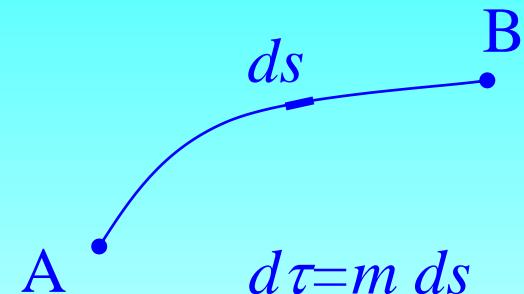
Fundamentals of geometrical optics

optical path, phase and intensity

Propagation of light in a medium:

Optical path:

$$\tau(AB) = \int_A^B m(s) ds$$



Complex refractive index:

$$\tilde{m} = m_r - im_i$$

- Phase difference : $\Delta\phi(AB) = k\tau_r = km_r \Delta S$
- Absorption: $I(B) = I(A) \exp(-2k\tau_i) = I(A) \exp(-2km_i \Delta S)$
- If m is constant: $\boxed{\Delta\phi = km_r \Delta S}$

$$\boxed{I = I_0 \exp(-2km_i \Delta S)}$$

Geometrical optics for light scattering

Can GO be applied to light scattering ?

YES to the simple shaped objects/particles:

- Homogeneous sphere,
- Homogeneous circular cylinder

Possible for objects of complex shape:

- Pure ray model, precision is very limited
- Ray model + electromagnetic integration, much better
- Ray model + wave properties → Vectorial Complex Ray Model
→ New concept, very precise and easy to use.

Geometrical optics for light scattering

Application of GO to light scattering

Light scattering by a sphere:

- Intensity :

$$\begin{aligned}\epsilon_X &= r_X^2 & p = 0 \\ \epsilon_X &= r_X^{2(p-1)}(1 - r_X^2)^2 & p \geq 1\end{aligned}$$

where p is the scattering order,
 X represents the polarization state
 $(\perp$ or \parallel), r_X Fresnel coefficients.

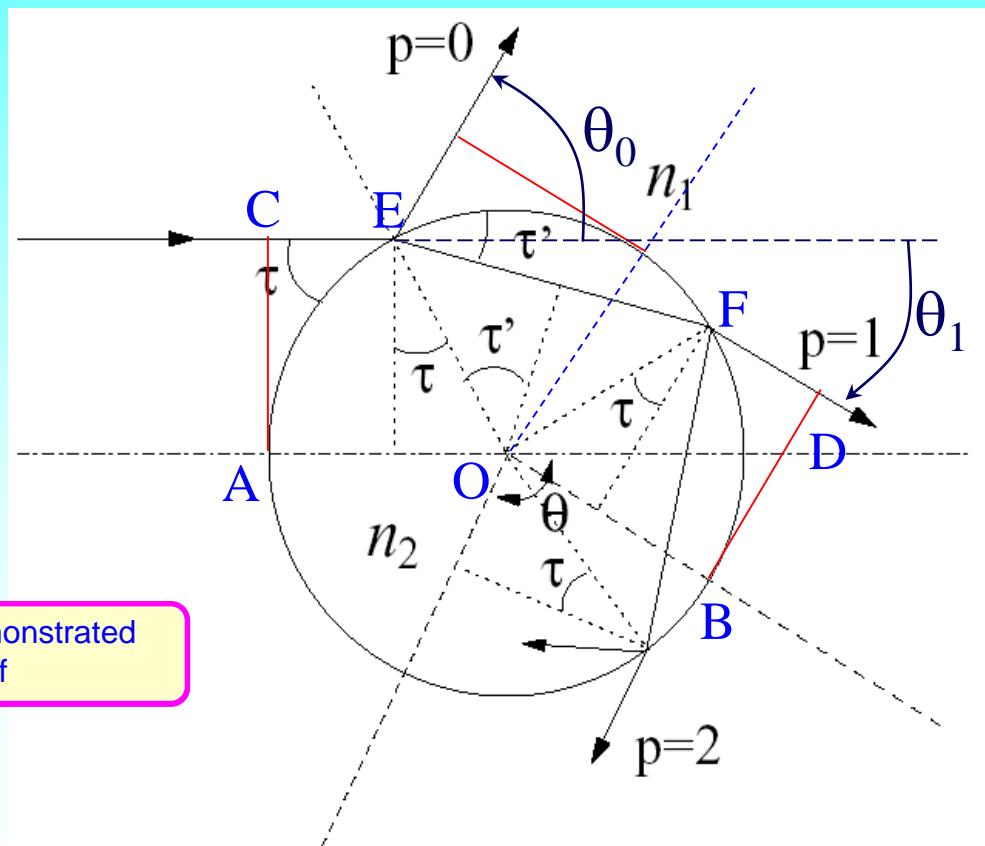
- Length of the path in the particle:

$$\Delta_p = 2pa \sin \tau'$$

- Deviation angle δ_e :

To be demonstrated
by yourself

$$\theta_p = 2\tau - 2p\tau'$$



Geometrical optics for light scattering

Application of GO to light scattering

➤ Phases:

- Phase difference due to the difference of the optical paths :

$$\Delta\phi = \frac{2\pi d}{\lambda} (\sin \tau - pm \sin \tau')$$

- Jump of phase due to reflection on the surface and the focal lines.

1) Reflection: phase of the complex number: r_x ,

2) Focal lines: phase jump $\pi/2$ at each focal line.

➤ Divergence factor :

$$D_p = \frac{\cos \theta_i \sin \tau}{\sin \theta_p \left| \frac{d\theta'_p}{d\tau} \right|}$$

$$I_{s,p} = \frac{I_0 \varepsilon_X dS_i}{dS_s} = \frac{I_0 \varepsilon_X a^2 \cos \tau \sin \tau d\tau d\phi}{r^2 \sin \theta_p d\theta_p d\phi} = \frac{a^2}{r^2} I_0 \varepsilon_X D_p$$

Summary

These equations are enough to calculate the intensity of each order and the total field.

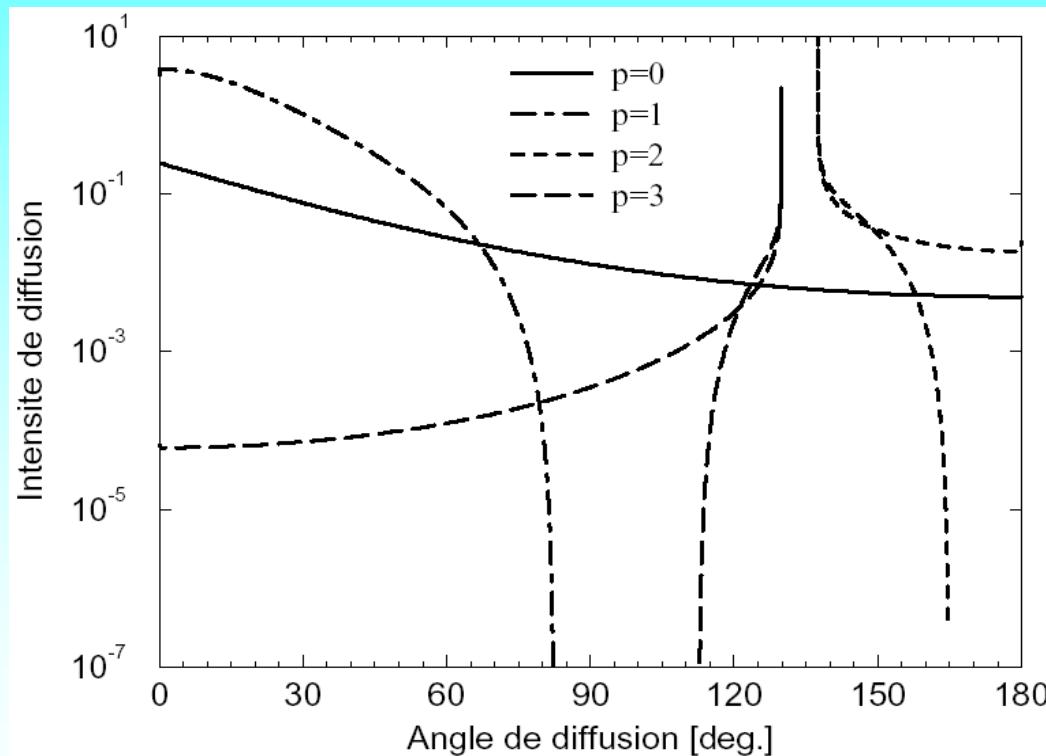
$$I_{s,t} = \sum_{p=0}^N \sqrt{I_{s,p}} e^{-ik\Delta\phi_p}$$

Geometrical optics for light scattering

Application of GO to light scattering

➤ Scattering diagram: **without** interference between different modes

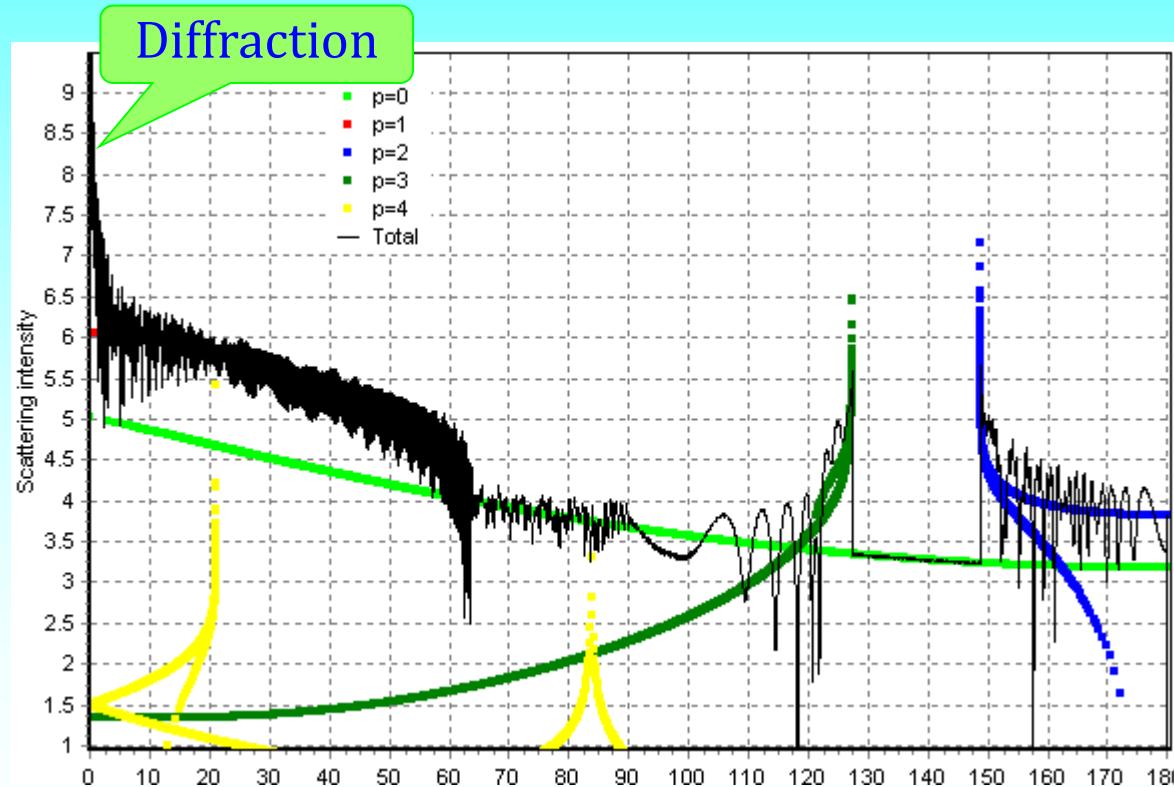
- Sphere of water
- Polarization \perp
- Intensity $\rightarrow \infty$ at rainbow angle



Geometrical optics for light scattering

Application of GO to light scattering

- Scattering diagram: With interference between different modes



Geometrical optics for light scattering

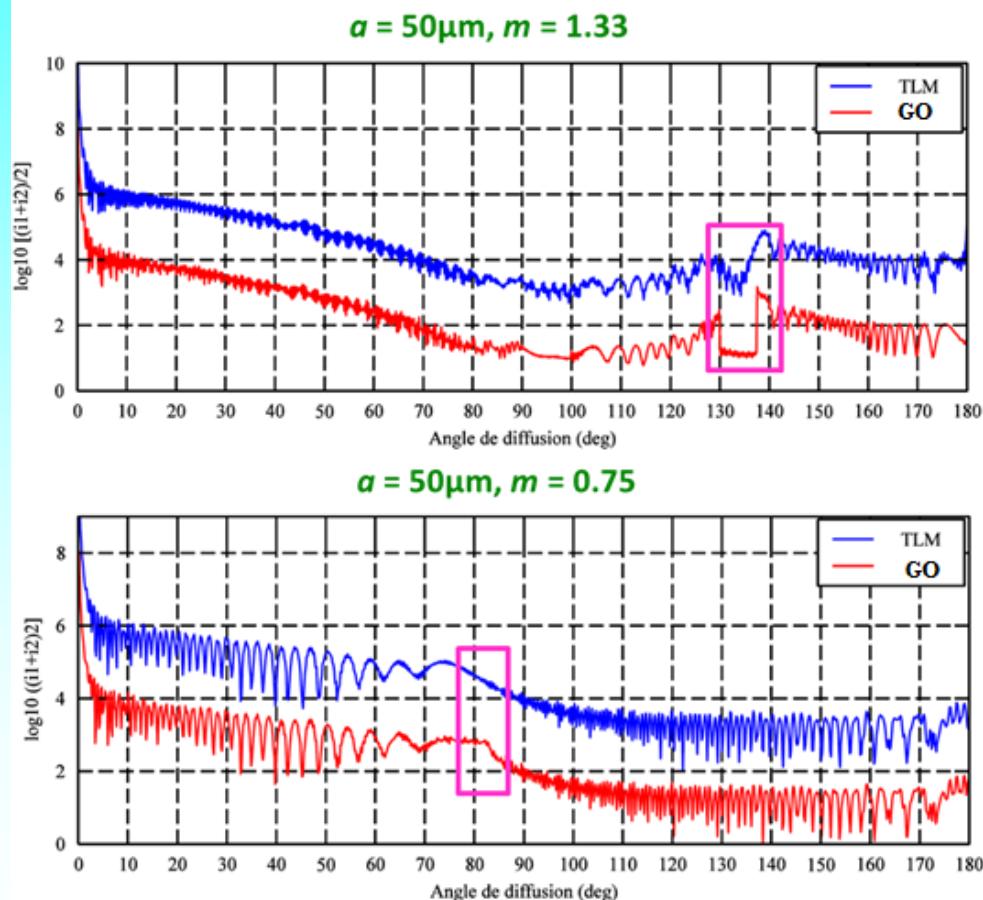
Application of GO to light scattering

➤ Scattering diagram:

Comparison with rigorous theory for a homogeneous sphere

Total intensity with interference:

A sphere of water or a air bubble in the water illuminated by a plane wave of wavelength of 0.6328 μm .



Geometrical optics for light scattering

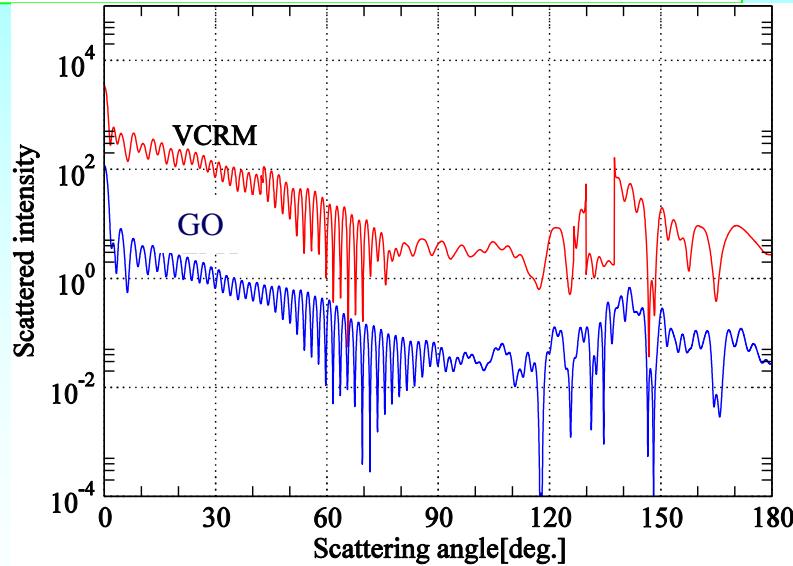
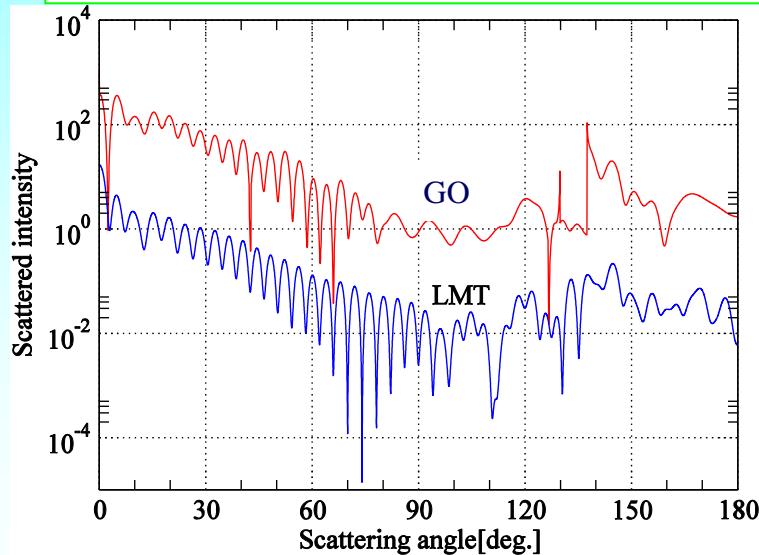
Application of GO to light scattering

➤ Scattering diagram: Comparison with rigorous theory for a cylinder

Total intensity with interference

A circular cylinder of water with a radius of 5 μm (left) or 10 μm (right) illuminated by a plane wave with a wavelength of 0.6328 μm .

Geometric optics still work very well for $a \sim 10\lambda$.



Geometrical optics for light scattering

Preliminary conclusions

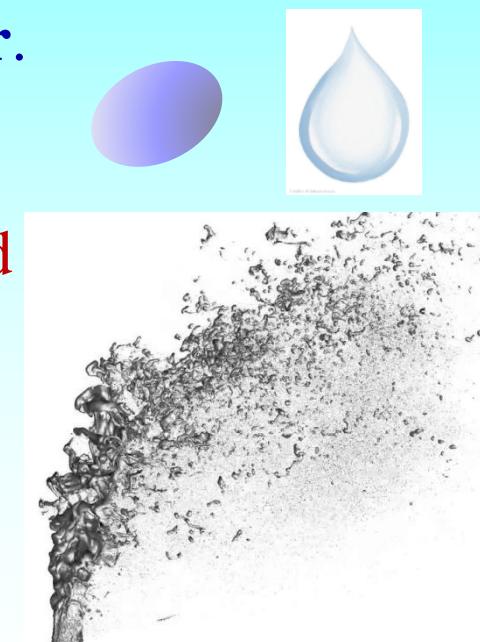
- By taking into account correctly the interferences, the Ray model can predict the scattering diagram in **ALL directions**.
- It can be applied to the scattering of **any shaped beam**.
- It works also for a circular infinite **cylinder**.

Limitations

NOT appropriate for a spheroid or an ellipsoid
NO for any irregular shaped particles.

The key problem is **the divergence factor**.

How to improve the model ?

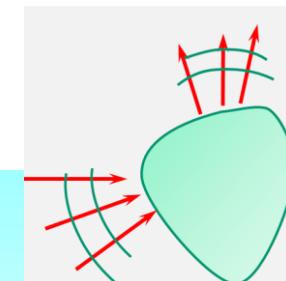


Vectorial Complex Ray Model

Possible candidate?

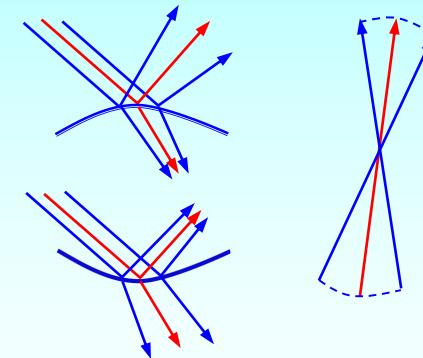
- * ~~Rigorous theories~~: the particle shape must correspond to a coordinates system
- * ~~Numerical methods~~: size limited, very time consuming
- ✓ Ray models: precision to be improved

Key problem: lack of wave properties



Our strategy: Extension of ray model

- Inclusion of wave front curvature
- Interference between all the rays.
- Diffraction.



Divergence/
Convergence

Phase in
focal lines

Vectorial Complex Ray Model

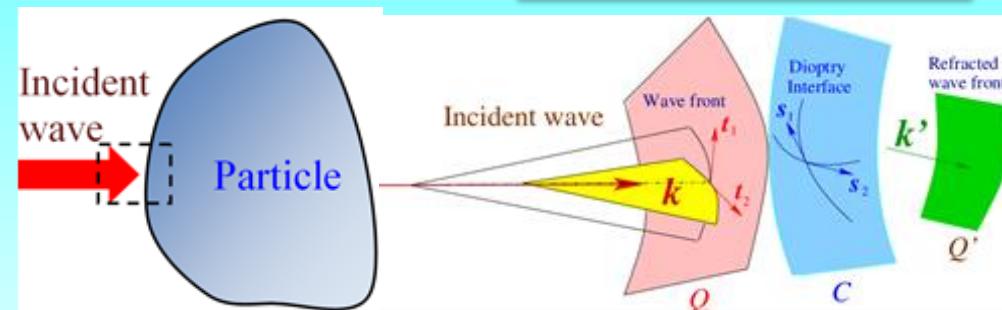
Vectorial Complex Ray Model

Geometrical optics + wave form

- **Vectorial Complex Ray Model – new**

- ✓ **5 properties of a ray:**

- Classical optics {
1. direction,
 2. amplitude,
 3. phase,
 4. polarization



For details:

- Ren et al, *Opt. Lett.* 36(3), 2011
- <http://www.amocops.eu>

important

New 5. Wave front curvature

- ✓ **Advantages:**

- Objects of any shape with smooth surface,
- Incident wave of any form,
- Sufficiently precise – scattering in all directions,
- All scattering properties of the objet.

Vectorial Complex Ray Model

Basic laws of geometrical optics

1. Snell-Descartes law:

Reflection: $i = i'$

Refraction: $n \sin i = n' \sin r$

2. Fresnel's Equations :

$$r_{\parallel} \equiv \frac{n_t \cos \theta_i - n_i \cos \theta_t}{n_i \cos \theta_t + n_t \cos \theta_i} = \frac{\tan(\theta_i - \theta_t)}{\tan(\theta_i + \theta_t)}$$

$$t_{\parallel} \equiv \frac{2n_i \cos \theta_i}{n_i \cos \theta_t + n_t \cos \theta_i} = \frac{2 \sin \theta_i \cos \theta_t}{\sin(\theta_i + \theta_t) \cos(\theta_i - \theta_t)}$$

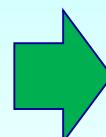
$$r_{\perp} \equiv \frac{n_i \cos \theta_i - n_t \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t} = -\frac{\sin(\theta_i - \theta_t)}{\sin(\theta_i + \theta_t)}$$

$$t_{\perp} \equiv \frac{2n_i \cos \theta_i}{n_i \cos \theta_i + n_t \cos \theta_t} = \frac{2 \sin \theta_i \cos \theta_t}{\sin(\theta_i + \theta_t)}$$

VCRM

The tangent component of wave vector is continuous

$$k_{\tau} = k'_{\tau}$$



$$r_{\perp} = \frac{k_{in} - k_{rn}}{k_{in} + k_{rn}}$$

$$t_{\perp} = \frac{2k_{in}}{k_{in} + k_{rn}}$$

$$r_{\parallel} = \frac{m^2 k_{in} - k_{rn}}{m^2 k_{in} + k_{rn}}$$

$$t_{\parallel} = \frac{2m k_{in}}{m^2 k_{in} + k_{rn}}$$

Vectorial Complex Ray Model

Summary of VCRM

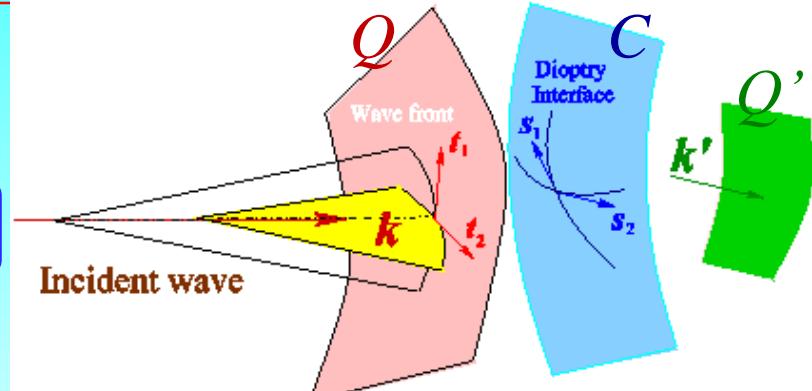
- Fundamental laws

1. Wave front equation:

$$(k_n^i - k_n^t)C = k' \Theta'^T Q' \Theta' - k \Theta^T Q \Theta$$

2. Law of Snell-Descartes in vectors:

$$k_\tau^i = k_\tau^t$$



- Amplitude:

$$A = \sqrt{D} \varepsilon$$

- Phase:

$$\Phi = \Phi_{inc} + \Phi_{fl} + \Phi_{path} + \Phi_\varepsilon$$

- Total field:

$$E = S_{diff} + \sum_{i=1}^N S_i$$

Divergence factor:

$$D = R_{11}^+ R_{21}^+ \frac{R_{12}^+ R_{22}^+}{R_{12}^- R_{22}^-} \dots \frac{R_{1q}^+ R_{2q}^+}{R_{1q}^- R_{2q}^-}$$

Fresnel coefficients: ε

$$\tilde{r}_\perp = \frac{k_n - \tilde{k}_n}{k_n + \tilde{k}_n} \quad \tilde{r}_\parallel = \frac{\tilde{m}^2 k_n - \tilde{k}_n}{\tilde{m}^2 k_n + \tilde{k}_n}$$

Vectorial Complex Ray Model

Special case of the wave front equation:

When the rays remain in the same plane – a main direction of the wave front and particle surface:

- Spherical particle
- Infinite cylinder at normal incidence
- Ellipsoidal particle in the symmetric plane.

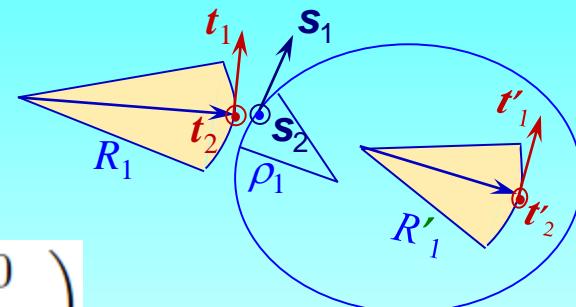
➤ Curvature matrix:

$$C = \begin{pmatrix} \frac{1}{\rho_1} & 0 \\ 0 & \frac{1}{\rho_2} \end{pmatrix}$$

$$\begin{aligned} Q &= \begin{pmatrix} \frac{1}{R_1} & 0 \\ 0 & \frac{1}{R_2} \end{pmatrix} \\ Q' &= \begin{pmatrix} \frac{1}{R'_1} & 0 \\ 0 & \frac{1}{R'_2} \end{pmatrix} \end{aligned}$$

➤ Wave front equation:

$$\begin{aligned} \frac{k'_n}{k'R'_1} &= \frac{k_n^2}{kR_1} + \frac{k'_n - k_n}{\rho_1} \\ \frac{k'}{R'_2} &= \frac{k}{R_2} + \frac{k'_n - k_n}{\rho_2} \end{aligned}$$



Rewrite the wave front equations in this simple form for the special case.

Vectorial Complex Ray Model

Applications to a sphere and a cylinder

➤ Sphere

- Reflection:

$$\dot{R}_1 = -\frac{a \cos \alpha}{2} \quad \dot{R}_2 = -\frac{1}{2a \cos \alpha} \rightarrow D = \frac{1}{4}$$

Derive these results from the wavefront equations (see below)

- Refraction $p=1$:

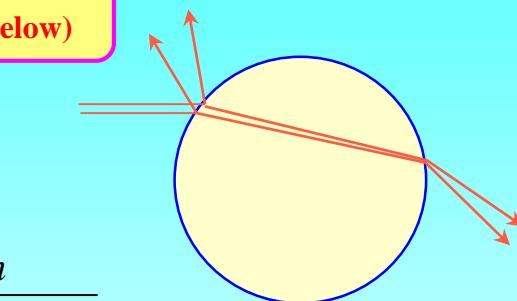
$$\text{After 1st refraction: } R'_{11} = -\frac{am \cos^2 \beta}{m \cos \beta - \cos \alpha} \quad R'_{12} = -\frac{am}{m \cos \beta - \cos \alpha}$$

After 2nd refraction:

$$R'_{21} = \frac{m \cos \beta - 2 \cos \alpha}{2(m \cos \beta - \cos \alpha)} \cos \alpha \quad R'_{22} = \frac{2a \cos \beta(m \cos \beta - \cos \alpha) - m}{2(m \cos \beta - \cos \alpha)(\sin \alpha \sin \beta - \sin \alpha \sin \beta)}$$

Divergence factor:

$$D = \frac{m \sin(2\alpha) \cos \beta}{4 \sin[2(\beta - \alpha)] (\cos \alpha - m \cos \beta)}$$



Check these by yourself.

Identical to the classical one.

➤ Cylinder: $R_2 = \infty$

- Reflection :

$$D = \frac{a \cos \alpha}{2}$$

- Refraction $p=1$:

$$D = \frac{m \cos \alpha \cos \beta}{2(\cos \alpha - m \cos \beta)}$$

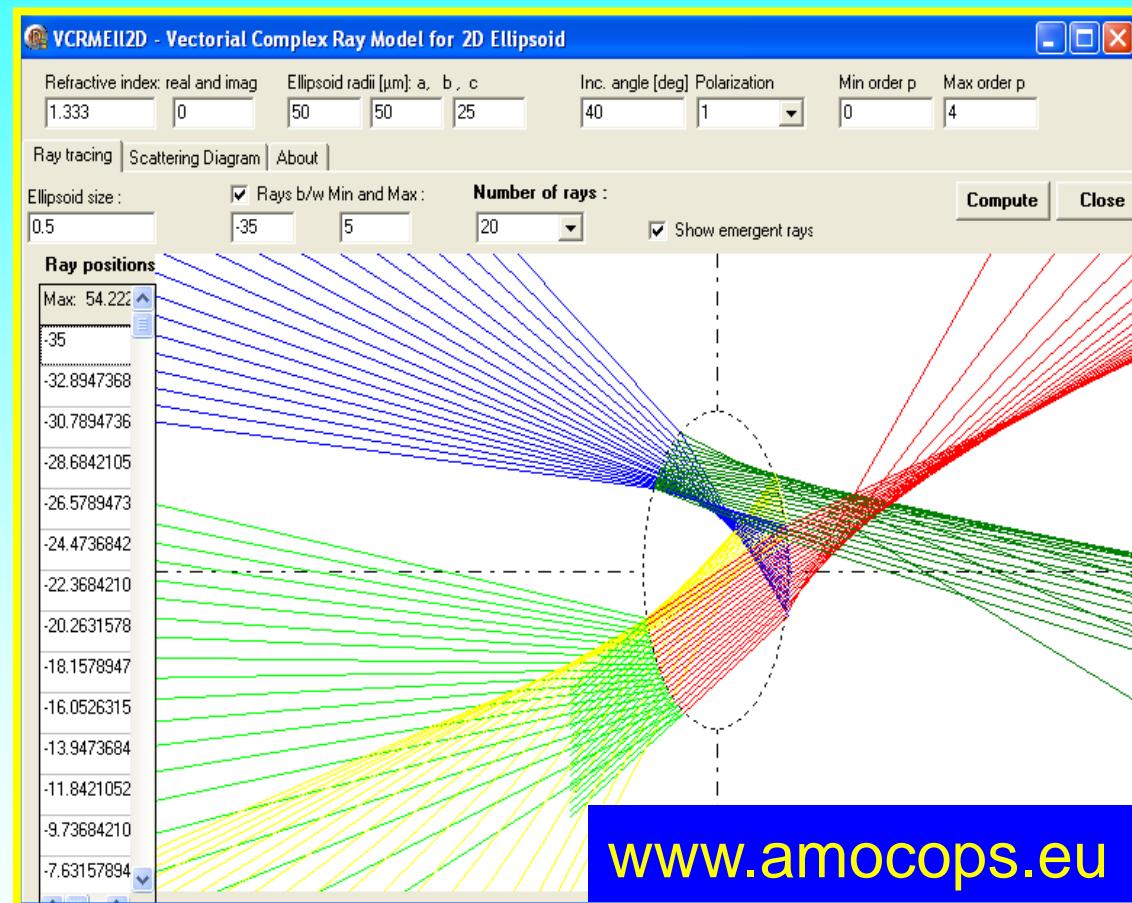
$$\begin{aligned} \frac{k'_n}{k' R'_1} &= \frac{k_n^2}{k R_1} + \frac{k'_n - k_n}{\rho_1} \\ \frac{k'}{R'_2} &= \frac{k}{R_2} + \frac{k'_n - k_n}{\rho_2} \end{aligned}$$

Vectorial Complex Ray Model

Software for an ellipsoid

Ray tracing Module

- Tracing of all the rays or a part of them
- Any number of rays
- Any position of rays

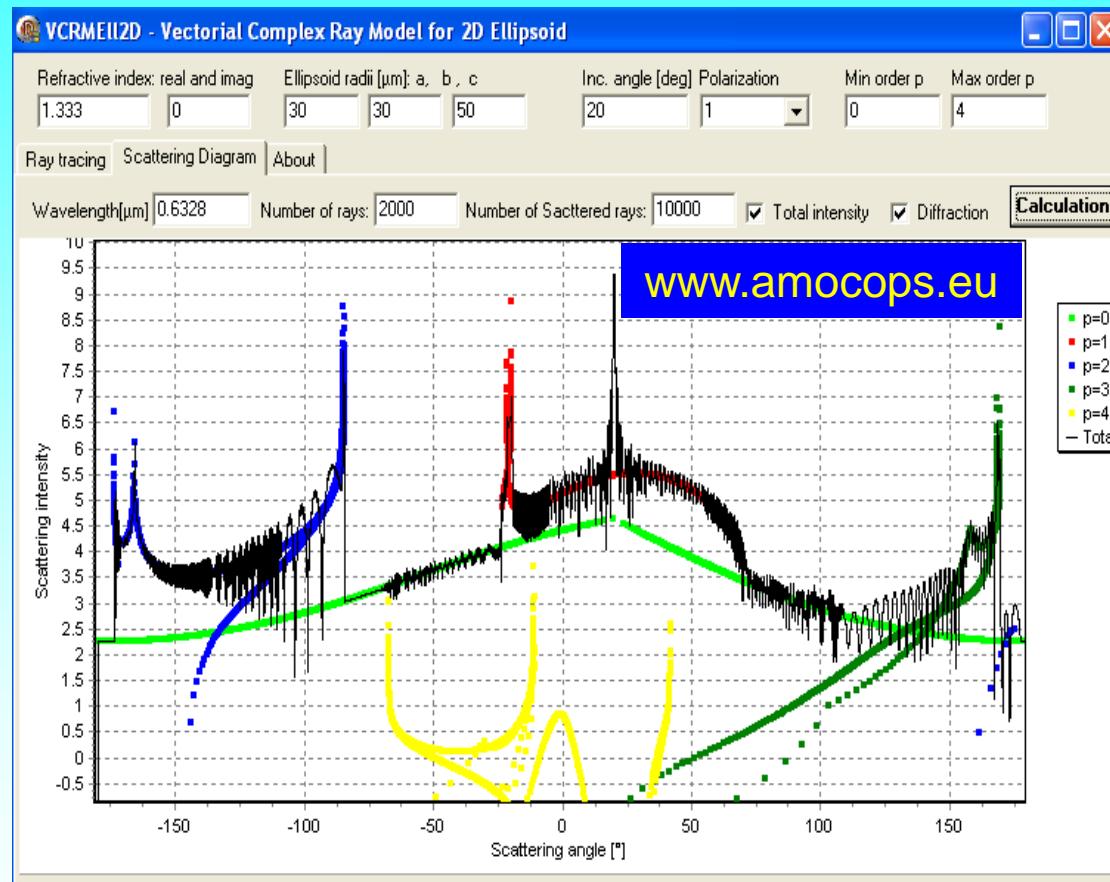


Vectorial Complex Ray Model

Software for an ellipsoid

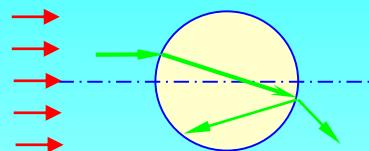
Scattering diagram module

- Intensity of each order
- total intensity of all the rays:
 - Interference
 - Diffraction

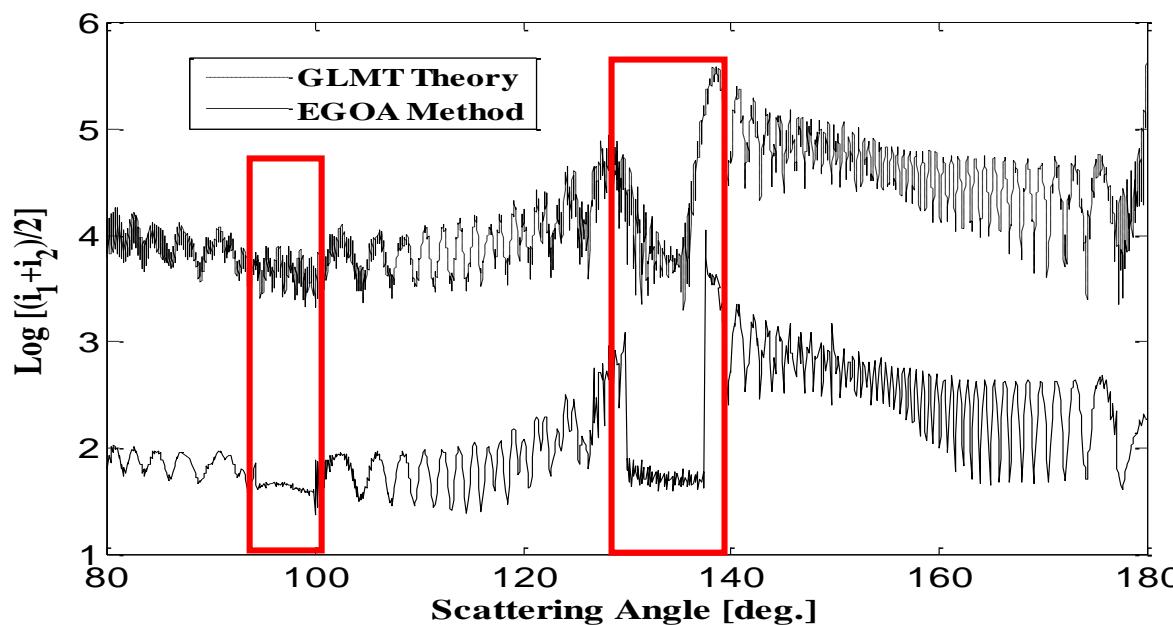


Vectorial Complex Ray Model

Comparison with Lorenz-Mie Theory

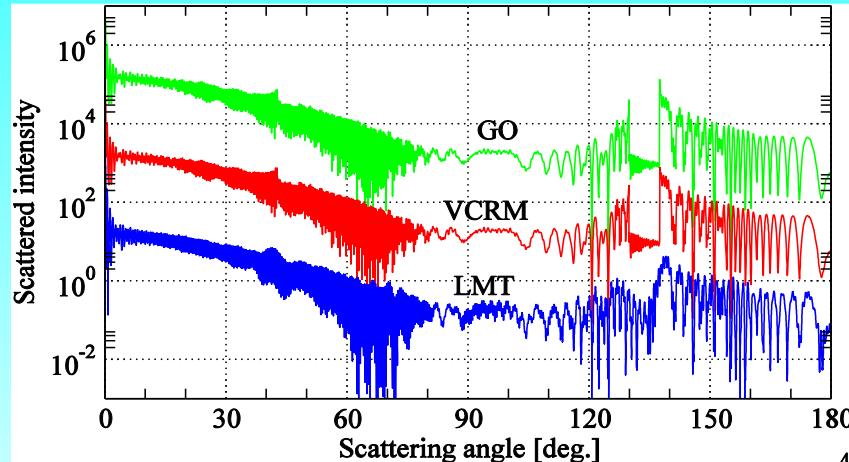


Sphere $a=100 \mu\text{m}$ **plane wave**
Discrepancy found near rainbow angles



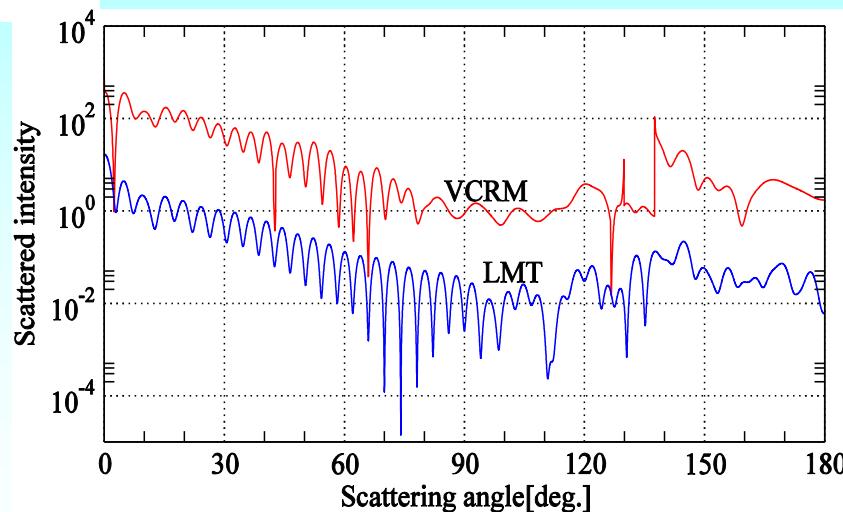
Vectorial Complex Ray Model

Comparison with Lorenz-Mie theory



← Scattering diagrams by LMT, GO and VCRM for an infinite circular cylinder :
 • refractive index: $m = 1.33$,
 • radius $a = 50 \mu\text{m}$
 • wavelength $\lambda = 0.6328 \mu\text{m}$.

→ Scattering diagrams by LMT, GO and VCRM for an infinite circular cylinder :
 • refractive index: $m = 1.33$,
 • radius $a = 5 \mu\text{m}$
 • wavelength $\lambda = 0.6328 \mu\text{m}$.

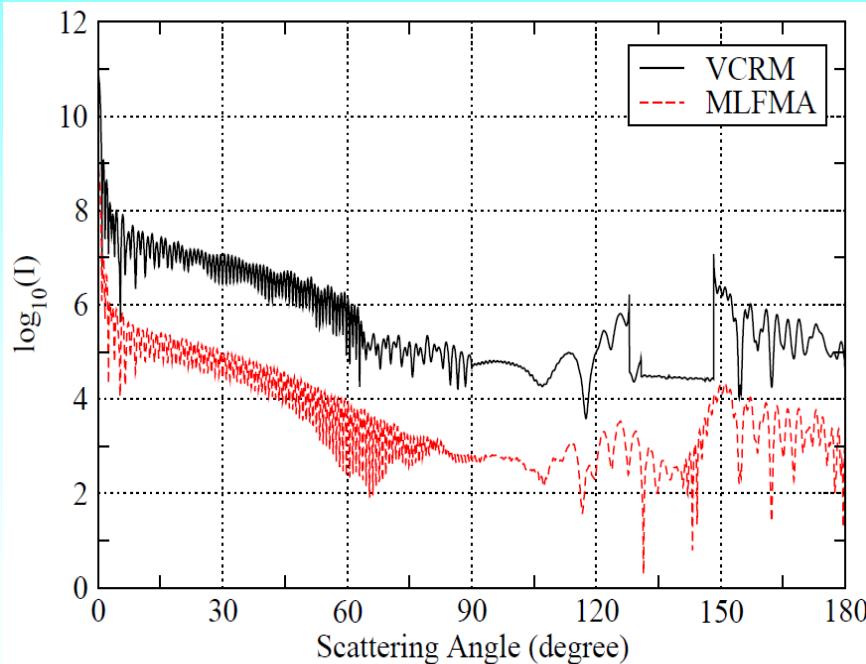


Vectorial Complex Ray Model

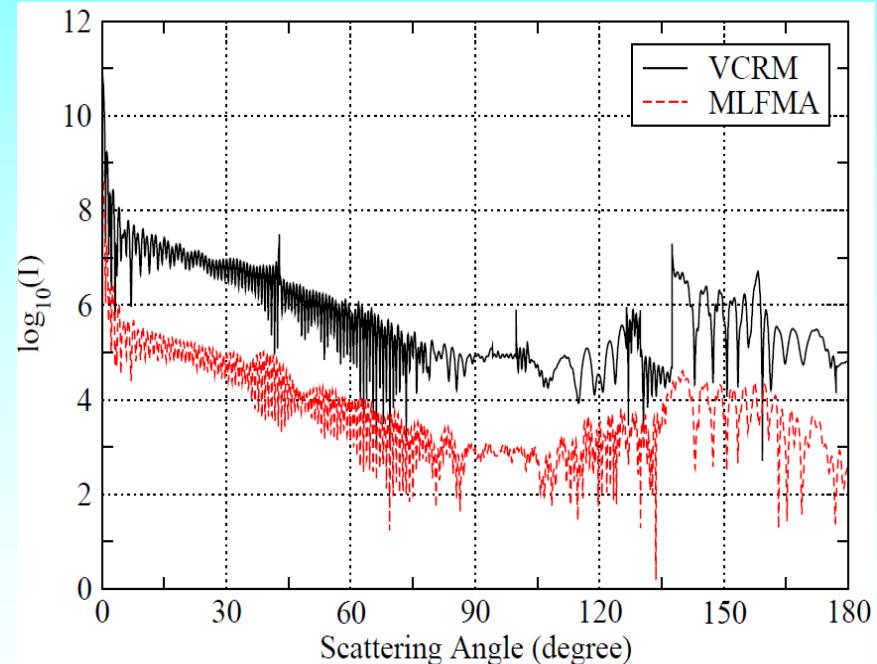
Comparison with exact numerical method MLFMA

A spheroidal water droplet ($m=1.33$) $a = 30 \mu\text{m}$ and a plane wave of wavelength $\lambda=0.785 \mu\text{m}$.

prolate: $a = b$, $c = 1.1a$



oblate: $a = 0.9b$, $c = a$



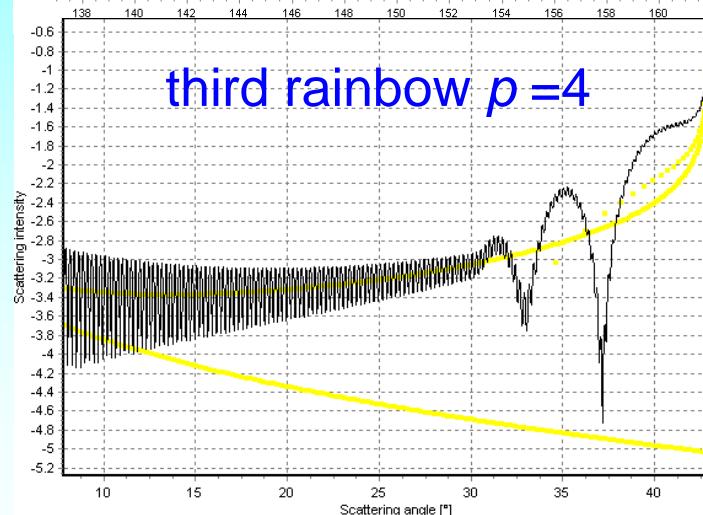
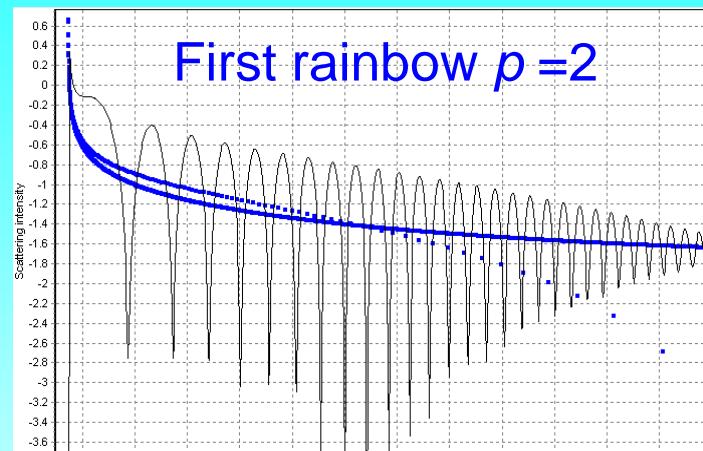
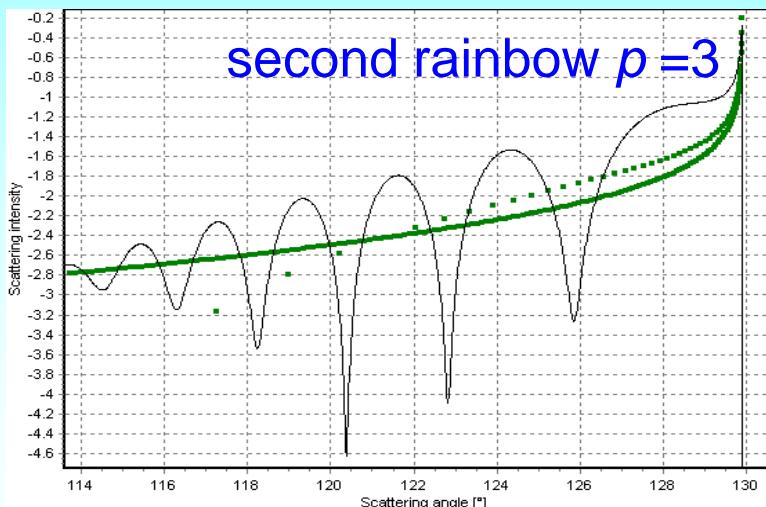
Vectorial Complex Ray Model

Interference near rainbow angles

Plane wave incident on a sphere:

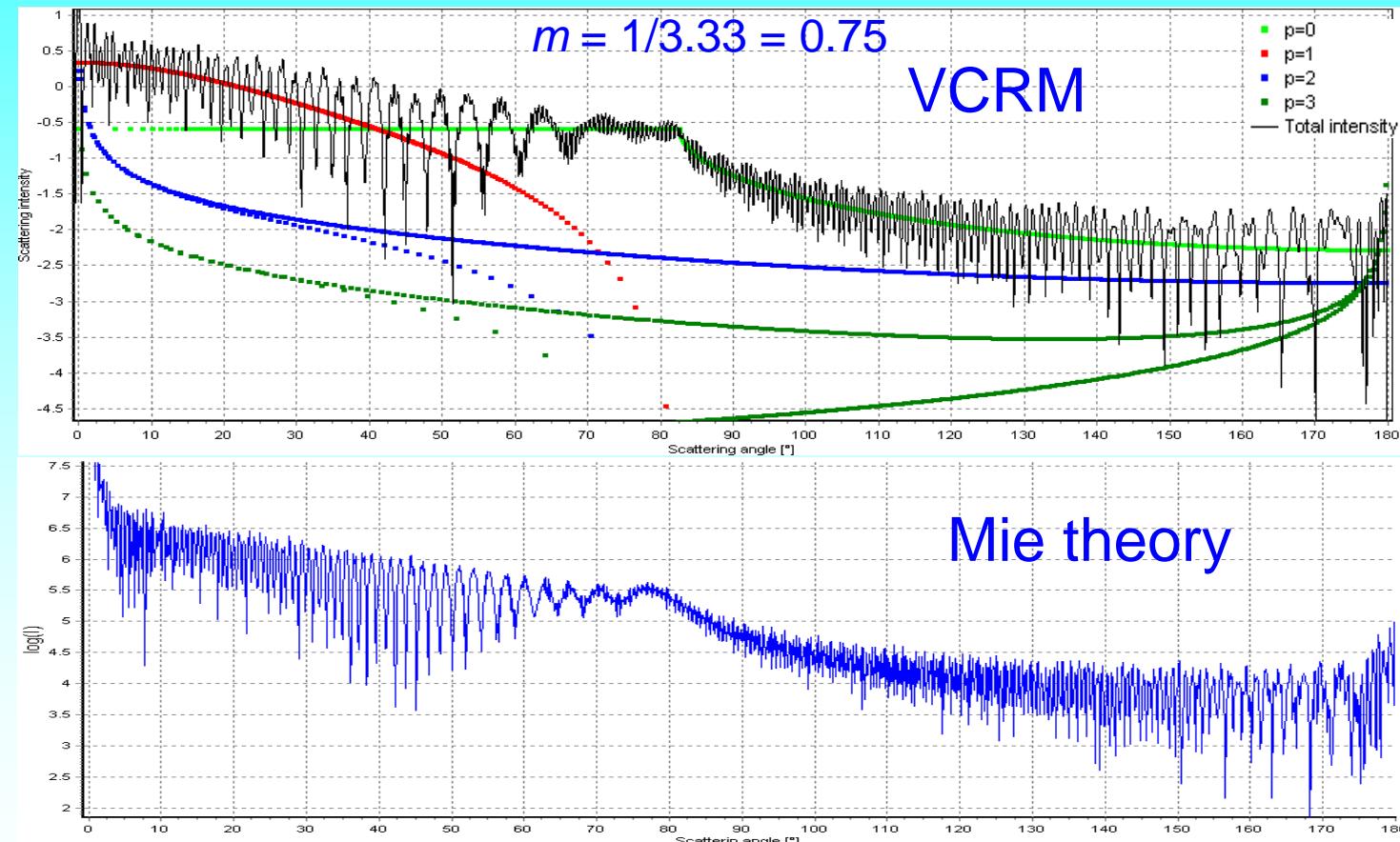
$\lambda = 0.6328 \mu\text{m}$,

$a = 100 \mu\text{m}$, $m = 1.33$



Vectorial Complex Ray Model

Comparison with Mie theory for a bubble



Ray Theory of Wave

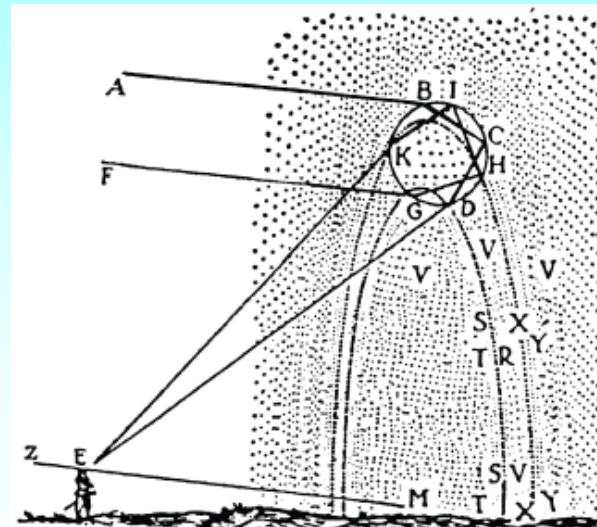
The rainbow is a marvel of nature.

L'arc-en-ciel est une merveille de la Nature si remarquable, et sa cause a été de tout temps si curieusement recherchée par les bons esprits, et si peu connue, que je ne saurais choisir de matière plus propre à faire voir comment, par la méthode dont je me sers, on peut venir à des connaissances que ceux dont nous avons les écrits n'ont point eues.

彩虹乃自然之奇观
自古引圣贤索其源



The Colors of bows are due to the refractive index.
But the intensity tends to infinity !!



Ray Theory of Wave

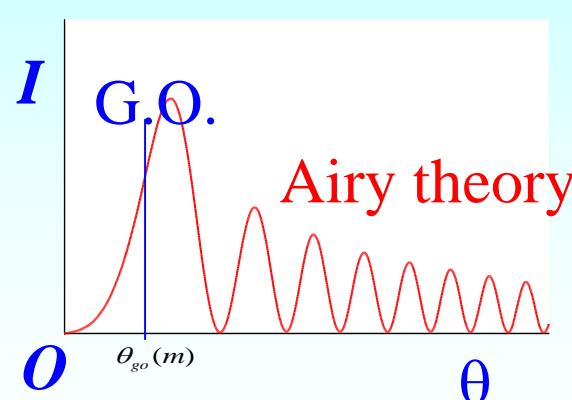
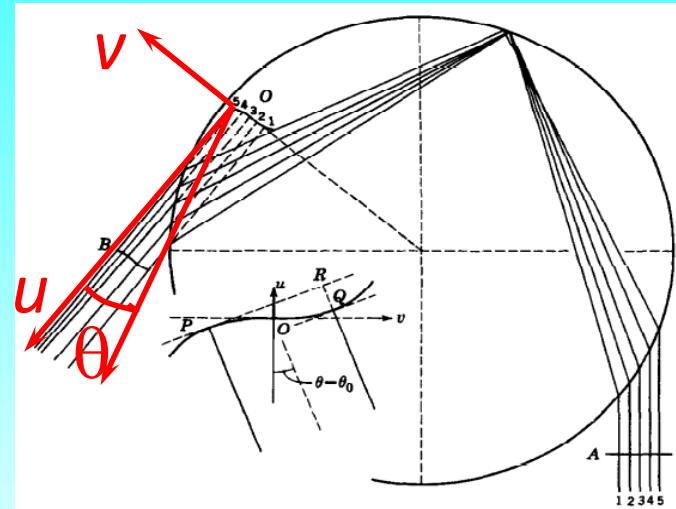
Airy theory (1838)



- Phase difference is:

$$\Delta\Psi = khv^3/3a^2$$
- Amplitude is constant for all emergent rays
- Amplitude of scattered field in θ direction:

$$\int_{-\infty}^{\infty} e^{-ikv(\theta - \theta_0) + ikhv^3/3a^2} dv$$



Airy theory predict a good profile,

Questions :

- The phase function found in $v \rightarrow 0$ but integration to infinity ?
- The amplitude is not constant near rainbow angles!
- It is only valid for a sphere.

Ray Theory of Wave

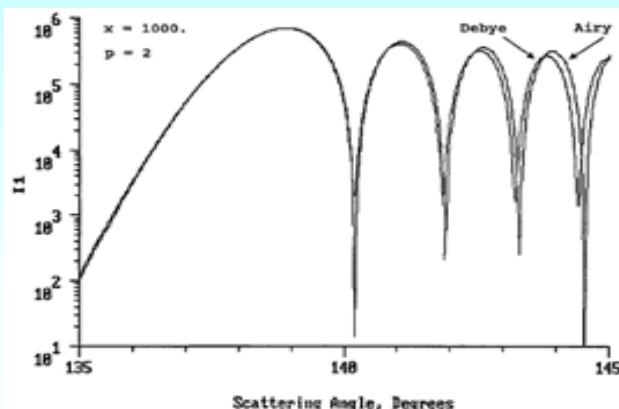
Airy theory in 1990s'

R. Wang and van de Hulst
Appl. Opt. 30(1):106, 1991

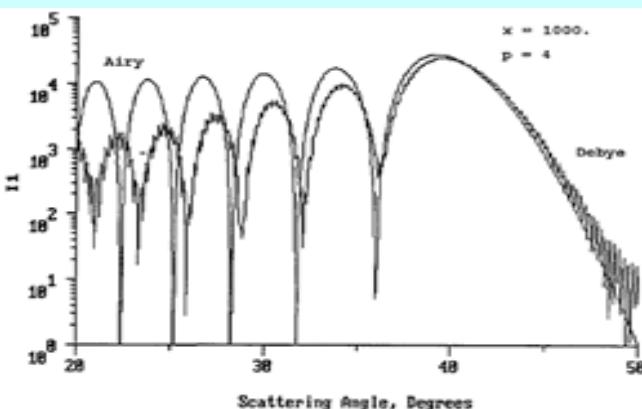
Can anything new be said about the rainbow? Yes. The insight that this phenomenon arises from the play of light in a single spherical drop is 7 centuries old, the full geometrical optics theory of Descartes 3-1/2 centuries, and its modification by Airy to take account of diffraction just 1-1/2 centuries. Exactly a century ago

► Wang, van de Hulst et Lock: same principle but with amplitude correction:

$$E_{\text{Airy}}^p(\theta) = x \left(\frac{2\pi \sin \theta_i^R}{\sin \theta^R} \right)^{1/2} \frac{x^{1/6}}{h^{1/3}} T^{21}(\theta_i^R) \times [R^{11}(\theta_i^R)]^{p-1} T^{12}(\theta_i^R) \times \\ \text{Ai}\left(\frac{-x^{2/3} \Delta}{h^{1/3}}\right) \exp(2\pi i L^R / \lambda) \exp\left[i(x\Delta) \left(\frac{p^2 - n^2}{p^2 - 1}\right)^{1/2}\right]$$



Airy theory
compared to
Debye theory
(rigorous)



Ray Theory of Wave

Airy theory in 1990s'

- Revisit of Airy theory to:
 1. understand the method,
 2. Extend to **non-spherical particle**.
- Theoretical demonstration for a sphere

- We know for each emergent ray:

phase: $\Phi = 2ka(\sin \tau - pm \sin \tau')$

deviation angle: $\theta = 2\tau - 2p\tau'$

- The derivatives: ($v \sim 0$):

$$\frac{du}{dv} \Big|_{\tau_0} = 0 \quad \frac{d^2u}{dv^2} \Big|_{\tau_0} = 0 \quad \frac{d^3\Delta\Phi}{d\tau^3} = -2kah \sin^3 \tau_0 = -\frac{2ka(p^2 - 1)}{p^2} \cos \tau_0$$

so
$$h = \frac{p^2 - 1}{p^2 \sin^2 \tau_0 \tan \tau_0} = \frac{(p^2 - 1)^2 \sqrt{p^2 - m^2}}{p^2(m^2 - 1)^{3/2}}$$

Ray Theory of Wave

Airy theory in 21st century

➤ **Phase:** $\Delta\Phi = \Phi(\mathbf{k}) - \Phi(\mathbf{k}_r) - k\overline{PR}$

with

$$\overline{OD} = r_0 \cdot \hat{k}_{r\perp}$$

$$\overline{OR} = r \cdot \hat{k}_\perp$$

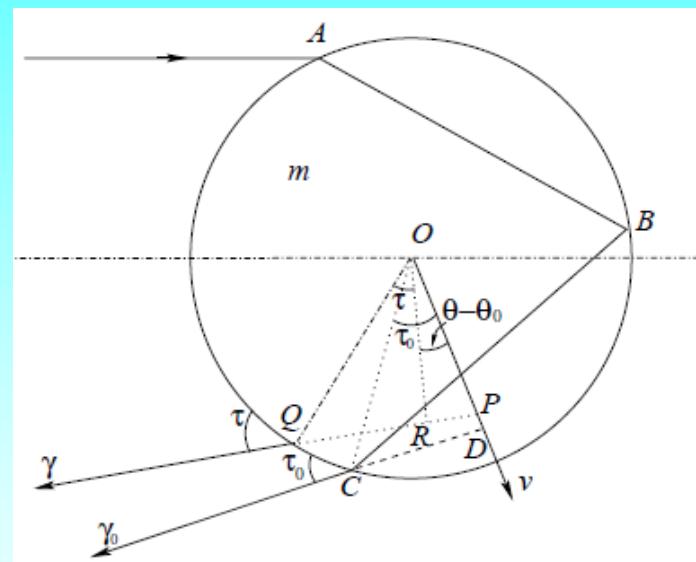
$$\overline{OP} = \frac{\overline{OR}}{\hat{k}_r \cdot \hat{k}} = \frac{\overline{OR}}{\hat{k}_{r+} \cdot \hat{k}_{+}}$$

$$\overline{PR} = \overline{OP}(\hat{k} \cdot \hat{k}_{r\perp})$$

$$v = \overline{OP} - \overline{OD}$$

➤ Amplitude

The amplitude of each ray is naturally calculated in VCRM.

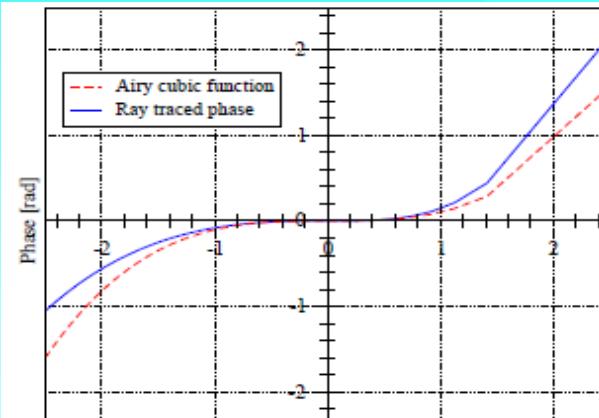
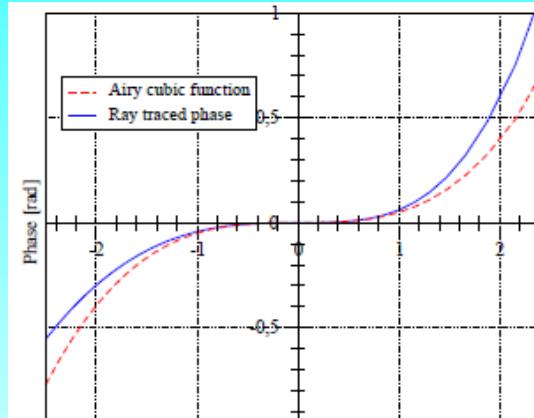


- All values are calculated **numerically**,
 - **Without any hypothesis/condition**.
 - The model is directly applicable to **any shaped particle**.

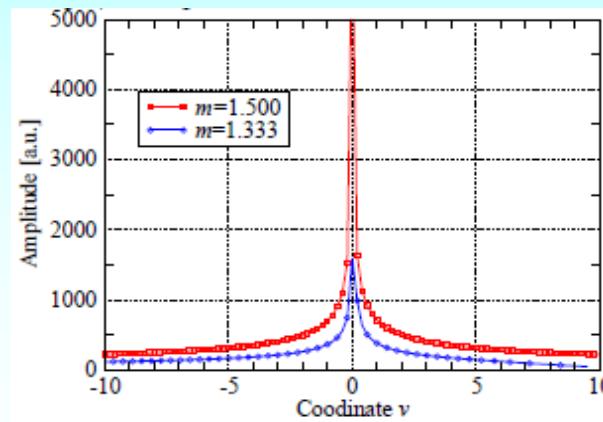
Ray Theory of Wave

Airy theory in 21st century

Phases calculated by Airy and VCRM

(a). $m = 1.333$ (b). $m = 1.5$

Amplitudes calculated by Airy and VCRM



Ray Theory of Wave

Toward the *Ray theory of wave*

VCRM can predict much better Airy structure than the Airy theory and can be applied directly to non-spherical particle.

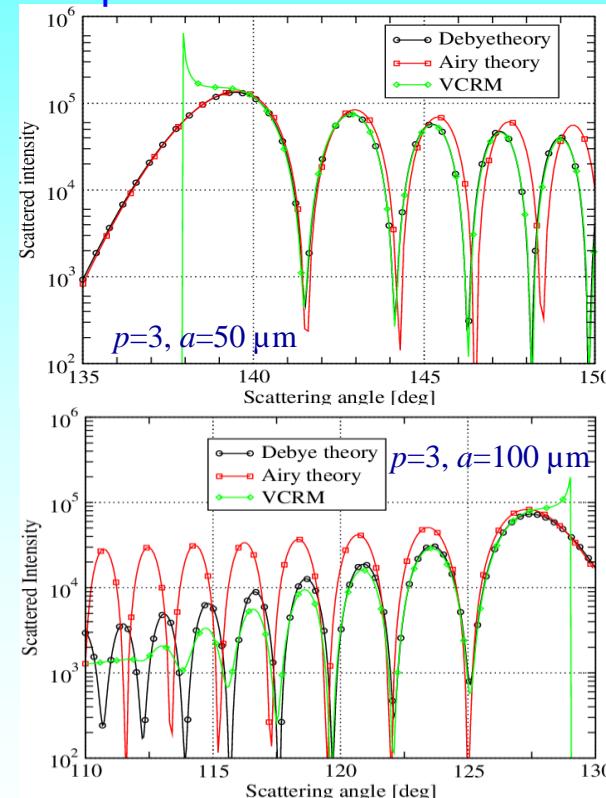
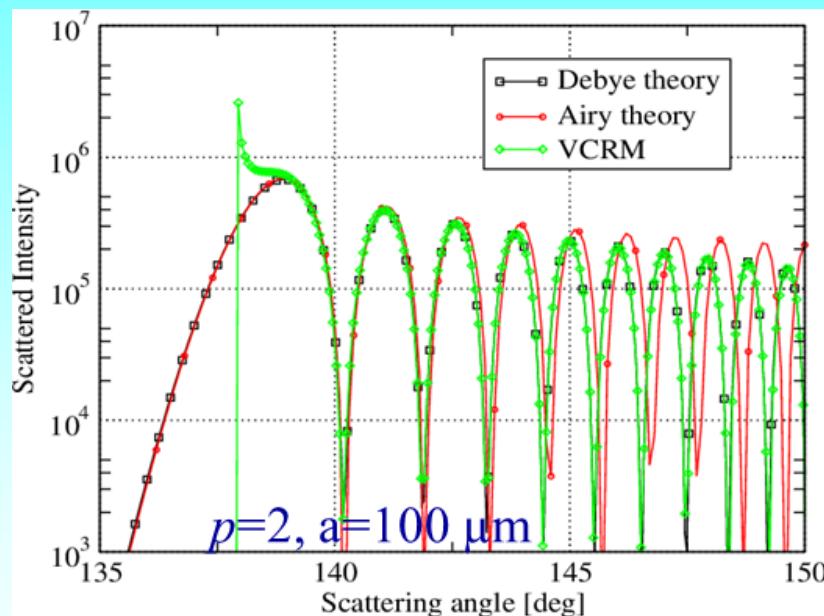
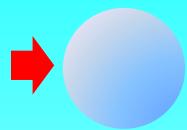
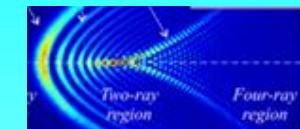


Fig. 3 Comparison of Airy structure calculated with the three methods.

State-of-art on Ray Theory of Wave



Motivations *An example*

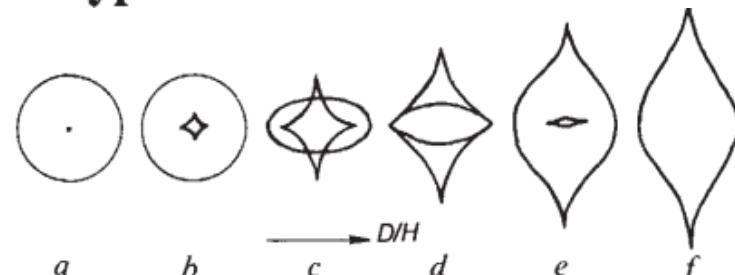


➤ Understanding of the natural phenomena:

- Fine structure of rainbow
- Optical caustics, cusp and catastrophes

[J. F. Nye, Nature \(London\) 312, 531 \(1984\)](#)

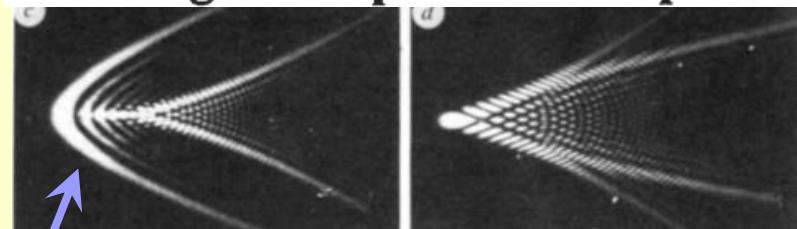
Rainbow scattering from spheroidal drops—an explanation of the hyperbolic umbilic foci



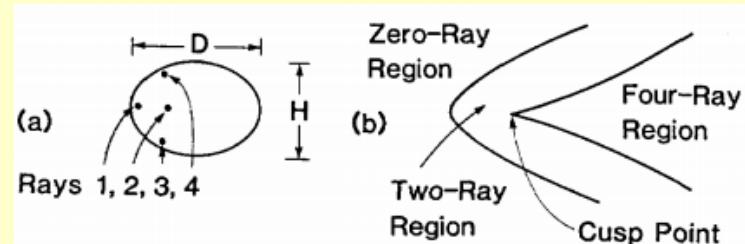
Can we predict the fine structure ?

[P. L. Marston, E. H. Trinh, Nature \(London\) 312, 529 \(1984\)](#)

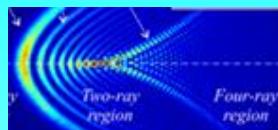
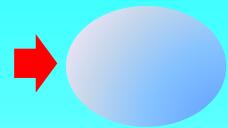
Hyperbolic umbilic diffraction catastrophe and rainbow scattering from spheroidal drops



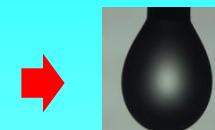
Marston, Opt. Lett. 1985



State-of-art on Ray Theory of Wave

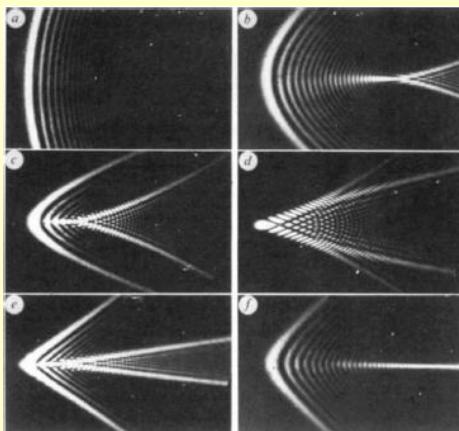


Motivations *An example*

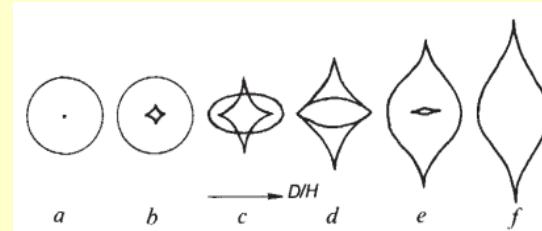


➤ Application to multiphase flow:

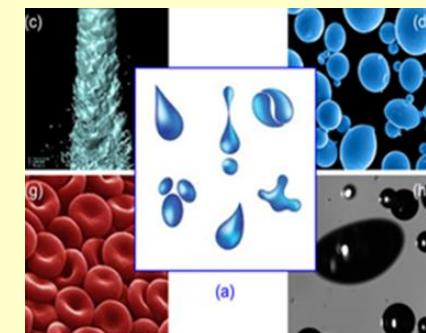
- Fine structure of rainbow
- Optical caustics, cusp and catastrophes



35 years ago
Marston → Nye



Today
1. Prediction of fine structure
2. Inversion



State-of-art on Ray Theory of Wave

Experimental validation

➤ Experimental set-up

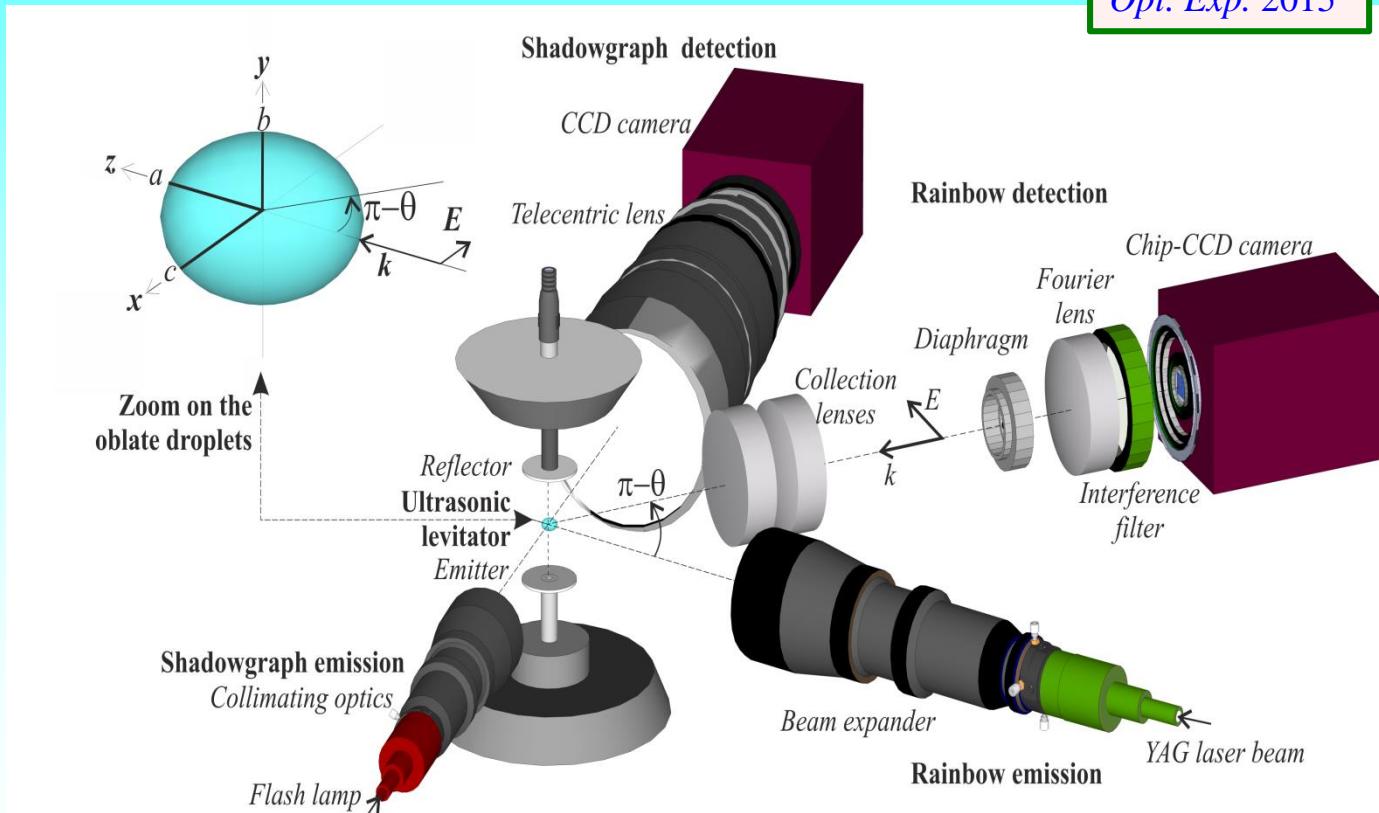
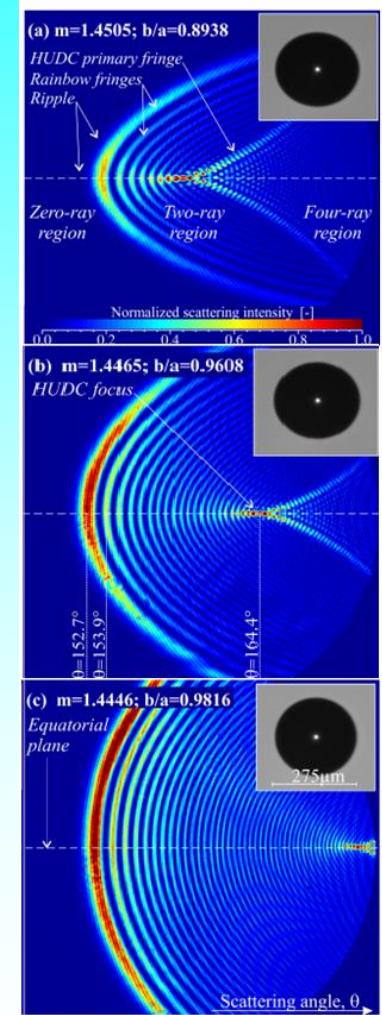


Figure 3 in
Onofri et al,
Opt. Exp. 2015



State-of-art on Ray Theory of Wave

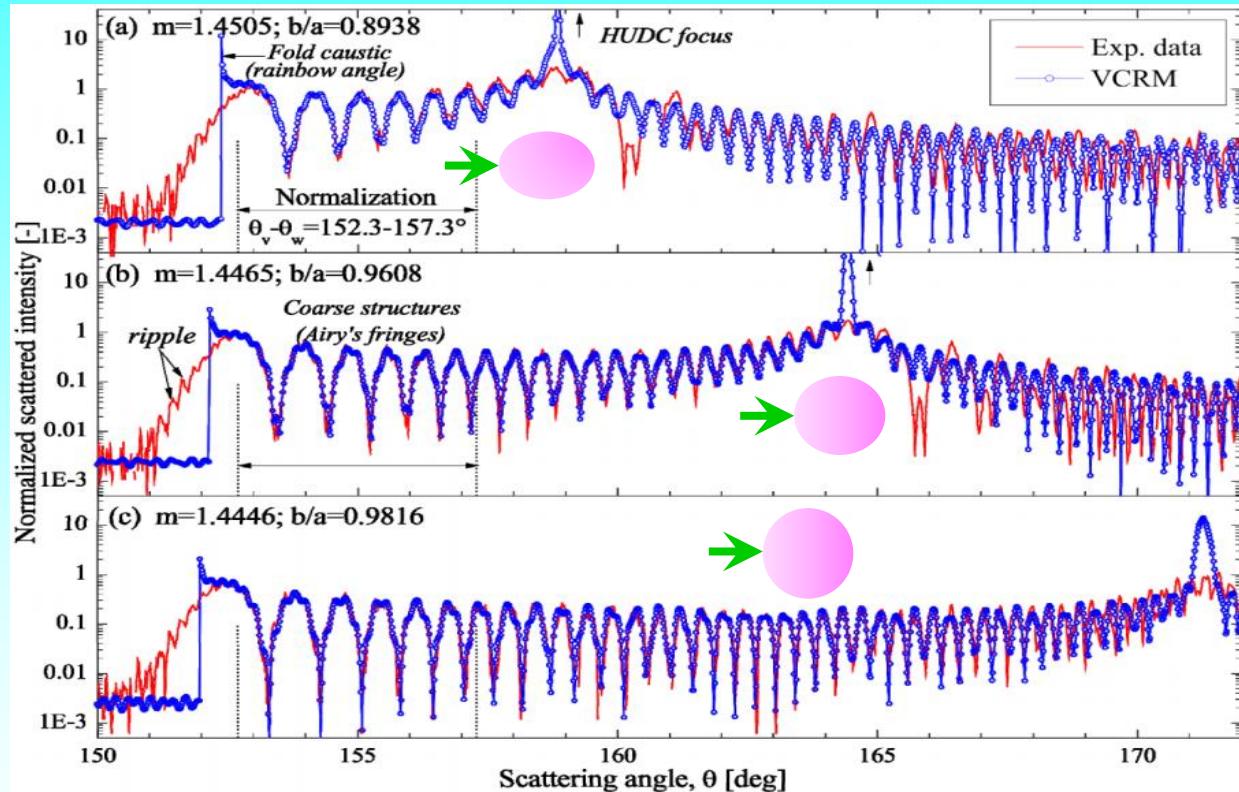
Experimental validation

- Comparison of the results

Figure 3 in
Onofri et al,
Opt. Exp. 2015

DEHS:
Di-Ethyl-Hexyl-
Sebacat

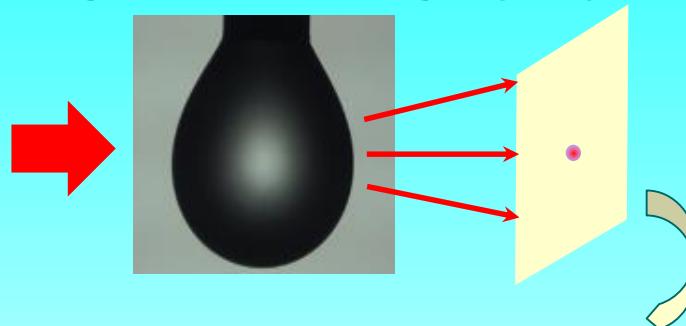
HUDC:
Hyperbolic
Umbilic
Diffraction
Catastrophe



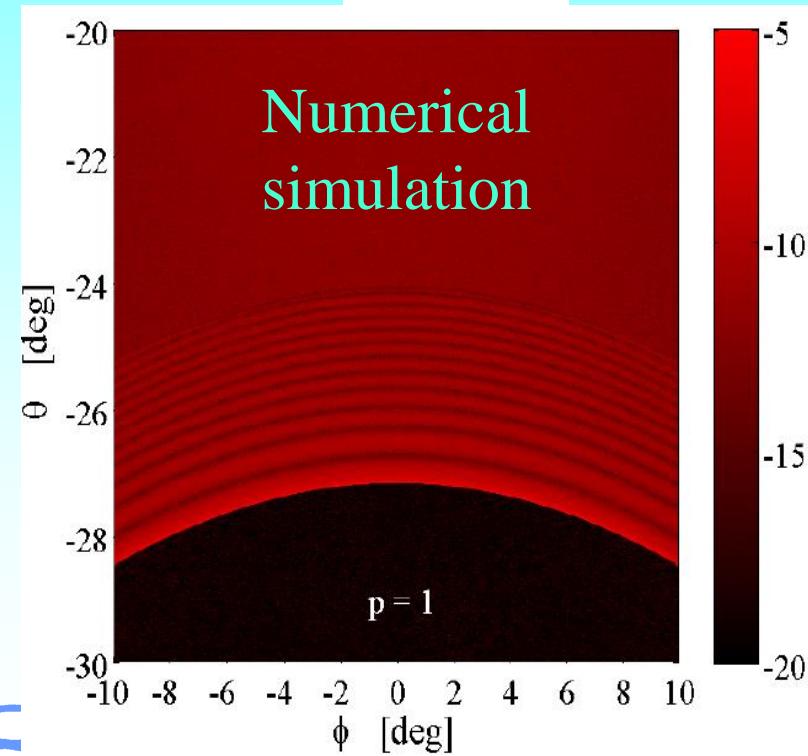
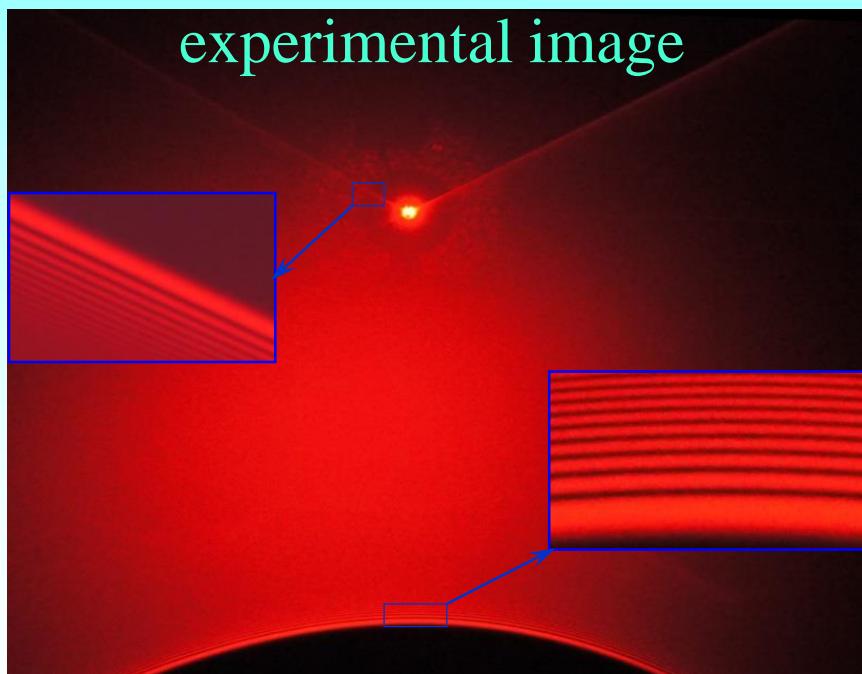
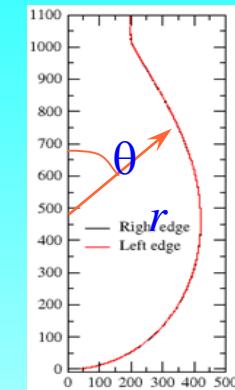
Comparison of VCRM and experimental normalized equatorial scattering diagrams for the droplets of 3 different aspect ratios. From (a) to (c), the droplet's aspect ratio b/a increases and refractive index decreases when the amplitude of the acoustic field is reduced.

State-of-art on Ray Theory of Wave

Light scattering by a pendant water drop

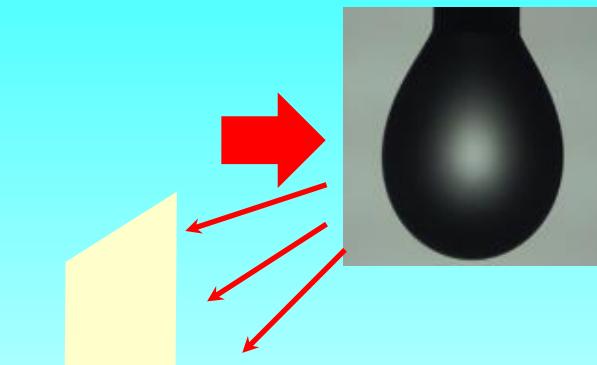


$$r(\theta) = a_0 + \sum_{i=2}^{10} a_i \theta^i$$

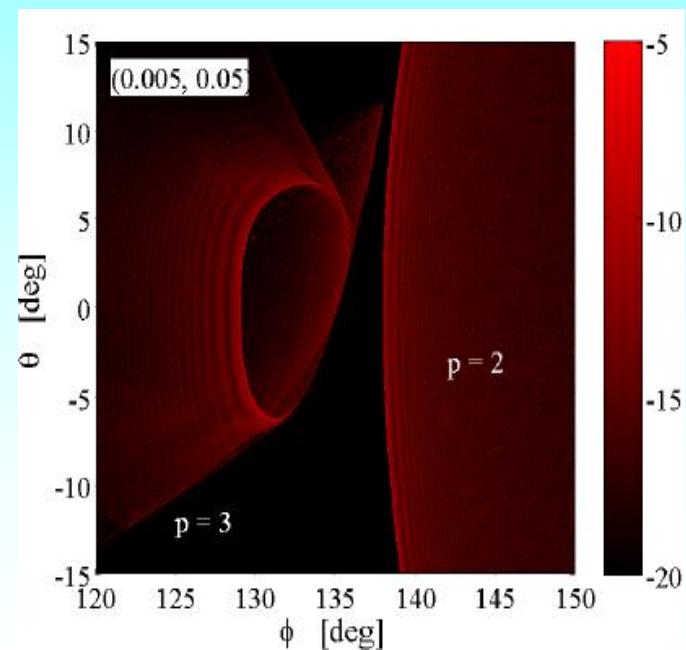
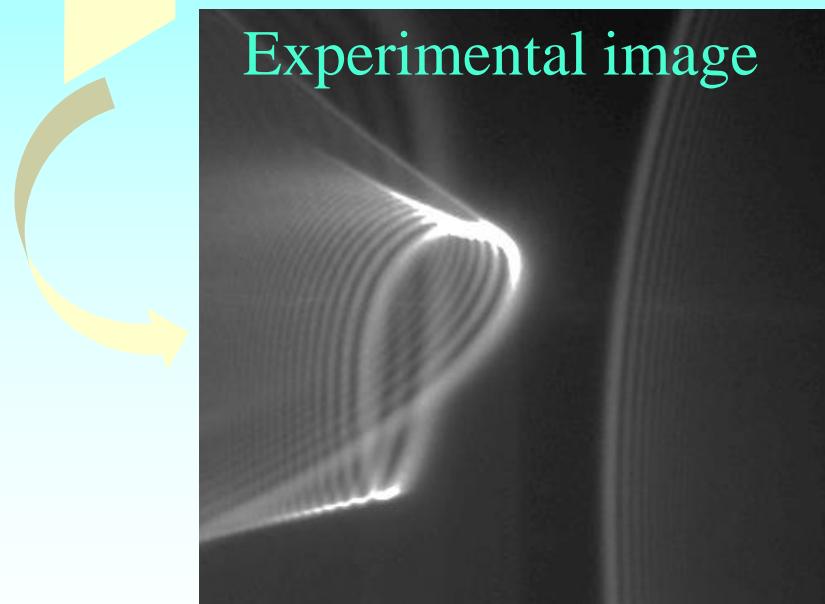
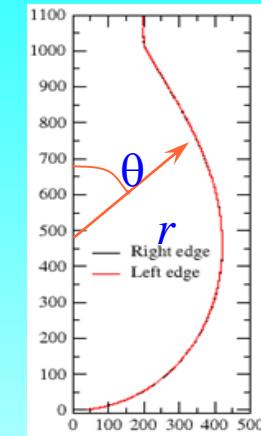


State-of-art on Ray Theory of Wave

Light scattering by a pendant water drop



$$r(\theta) = a_0 + \sum_{i=2}^{10} a_i \theta^i$$



State-of-art on Ray Theory of Wave

► Shaped beams in VCRM

- Electric field of a wave:

$$\tilde{E}(x, y, z) = E(x, y, z) e^{i\varphi(x, y, z)}$$

Where $E(x, y, z)$ is the amplitude and $\varphi(x, y, z)$ the phase.

- Propagation direction

$$\hat{k} = \frac{\nabla \varphi(x, y, z)}{\|\nabla \varphi(x, y, z)\|}$$

- Curvature matrix of the wave front surface:

The wave front surface being $\varphi(x, y, z) = \text{Const.}$ the principal directions and the principal curvatures are determined by the *generalized eigenvalues problem*:

$$\begin{vmatrix} L - \kappa E & M - \kappa F \\ M - \kappa F & N - \kappa G \end{vmatrix} = 0$$

E, F, G and L, M, N are respectively parameters of the first and the second fundamental forms of $\varphi = C$.

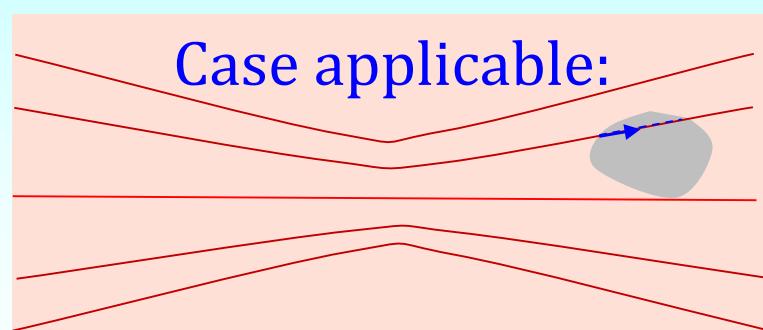
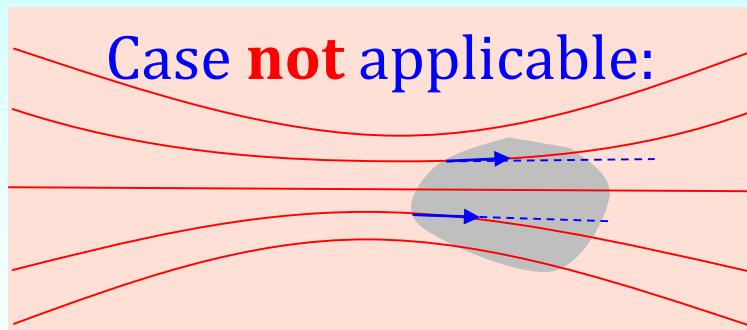
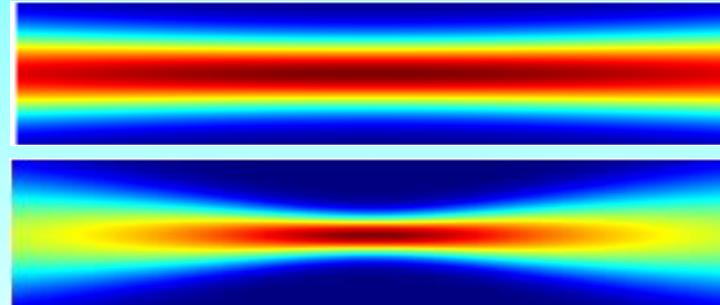
State-of-art on Ray Theory of Wave

➤ Shaped beams in VCRM

Light beam in VCRM

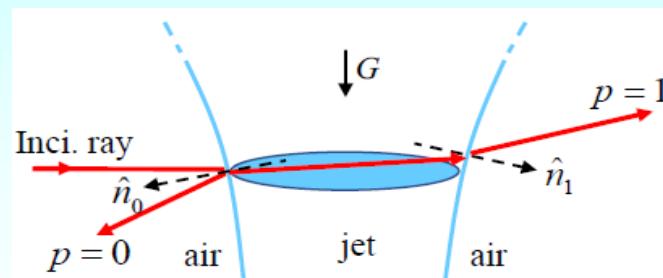
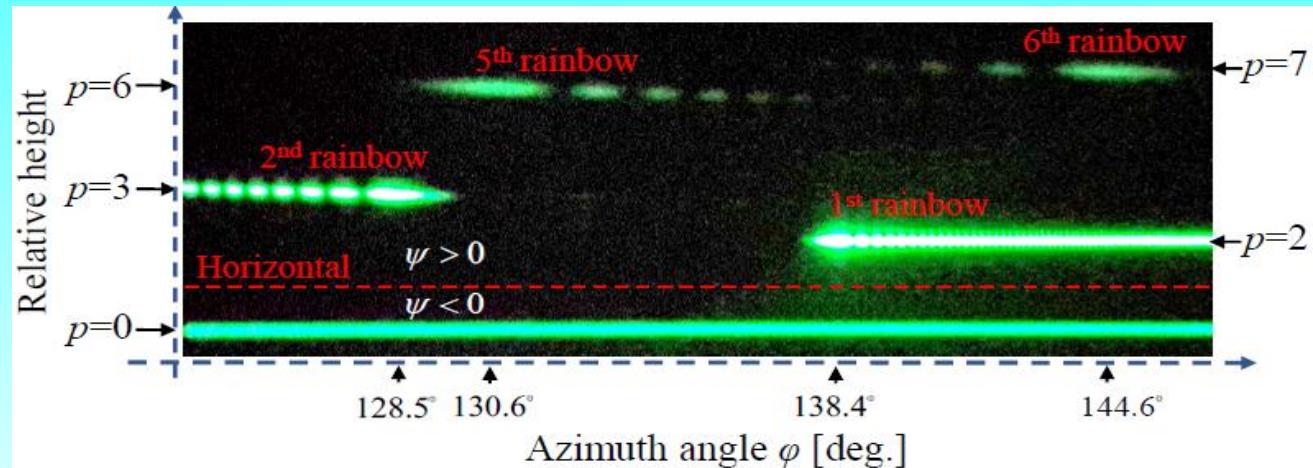
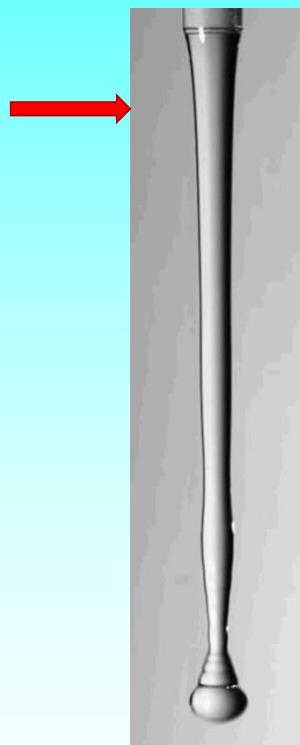
1. *The object is large compared to the wavelength.*
2. The rays propagate **rectilinearly / straightly** in a homogeneous medium, at least approximately in the scale of the object.

- No problem for collimated beams
- For focused beam,
The object should be small and far from the beam waist.



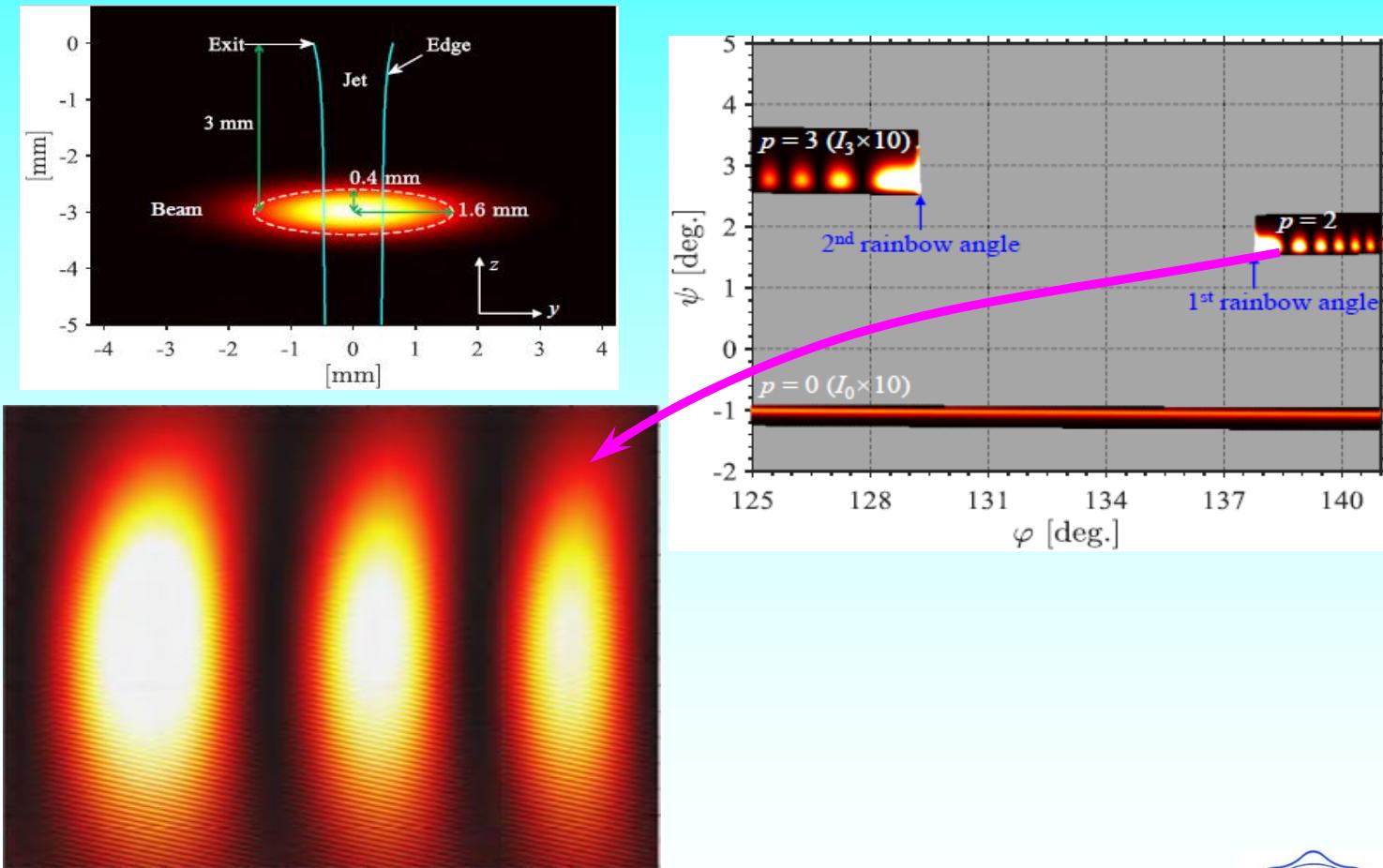
State-of-art on Ray Theory of Wave

- Scattering of an elliptical Gaussian beam by a real liquid jet
Experimental scattering pattern



State-of-art on Ray Theory of Wave

- Scattering of an elliptical Gaussian beam by a real liquid jet
Simulation with VCRM



Conclusions

- **Conclusions**

- The Vectorial Complex Ray model (VCRM) is established.
- VCRM has been validated :
 - theoretically for the spherical and cylindrical particle.
 - numerically for the scattering of a ellipsoid.
 - experimentally for the scattering of an oblate droplet, a real liquid jet and a pendant drop.
- A software for 2D ellipsoid is realized.
- VCRM can be applied to any shaped beams,
- First step to the Ray Theory of Wave achieved.

- **Application fields:**

- Metrology of non-spherical particles in multiphase flow,
- Prediction of radiation force, torque, and stress,
- Application to freeform optical systems,
- Diagnostics for micro-fluidics.

Thank you for your attention

